



Galaxy Power Spectrum and Biasing Results from the LOFAR Two-meter Sky Survey (First Data Release)

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Received 2021 September 25; revised 2022 February 18; accepted 2022 February 18; published 2022 March 25

Abstract

The Low-Frequency Array (LOFAR) Two-meter Sky Survey (LoTSS) is an ongoing survey aiming to observe the entire northern sky, providing an excellent opportunity to study the distribution and evolution of the large-scale structure of the universe. The source catalog from the public LoTSS first data release (DR1) covers 1% of the sky, and shows correlated noise or fluctuations of the flux density calibration on scales of a few degrees. We explore the LoTSS DR1 to understand the survey systematics and data quality of this first data release. We produce catalog mocks to estimate uncertainties, and measure the angular clustering statistics of LoTSS galaxies, which fit the Lambda cold dark matter cosmology reasonably well. We employ a Markov Chain Monte Carlo–based Bayesian analysis to recover the best galaxy biasing scheme and multicomponent source fraction for LoTSS DR1 above 1 mJy assuming different possible redshift templates. After masking some noisy and uneven patches and with suitable flux density cuts, the LOFAR survey appears qualified for large-scale cosmological studies. The upcoming data releases from LOFAR are expected to be deeper and wider, and will therefore provide improved cosmological measurements.

Unified Astronomy Thesaurus concepts: [Cosmology \(343\)](#); [Large-scale structure of the universe \(902\)](#); [Radio galaxies \(1343\)](#)

1. Introduction

Our present understanding of the origin and evolution of the universe is based on the Lambda cold dark matter (Λ CDM) cosmology. In this model, the matter density of the universe is dominated by cold dark matter, whose gravitational evolution results in a population of virialized dark matter halos of different masses (Press & Schechter 1974; Scherrer & Bertschinger 1991; Navarro et al. 1997; Ma & Fry 2000; Seljak 2000; Scoccimarro et al. 2001; Cooray & Sheth 2002). The formation of galaxies occurs inside these dark matter halos, and the host halo mass and its evolution correlates with the evolution and type (Girelli et al. 2020) of galaxy residing inside. In general, radio-loud galaxies' active galactic nuclei (AGNs) are found to reside in more massive halos than optical AGNs (Wilman et al. 2008; Mandelbaum et al. 2009). That being so, the optical and radio observations sample quite a different set of galaxies (halos). Most galaxies which are bright at optical wavelengths are undetectable at radio wavelengths, and strong radio sources are often optically faint or invisible. The radio surveys sample the higher end of the halo mass range as compared with optical observations, and thus complement existing and upcoming visible/IR galaxy surveys (Kauffmann et al. 2003; Best et al. 2005, 2014; Mandelbaum et al. 2009; Nusser & Tiwari 2015; Hale et al. 2018; Krumpel et al. 2018; Alonso et al. 2021; Lan & Prochaska 2021; Wolf et al. 2021).

The international Low-Frequency Array (LOFAR; van Haarlem et al. 2013) is a new-generation radio interferometer constructed in the northern Netherlands and across Europe, offering an increase in

radio survey speed with unparalleled sensitivity and angular resolution in the low-frequency radio regime (Röttgering et al. 2003, 2005; van Haarlem 2005; Falcke et al. 2007). The LOFAR Two-meter Sky Survey (LoTSS) is an ongoing deep 120–168 MHz imaging survey being carried out using LOFAR high-band antenna observations across the whole Northern Hemisphere (Shimwell et al. 2017, 2019). The LoTSS aims to explore cosmic large-scale structure, galaxies, clusters of galaxies, and the formation and evolution of massive black holes. The LoTSS survey will observe millions of radio AGNs, along with a significant number of star-forming galaxies (SFGs) out to redshift $z \sim 6$, allowing detailed studies of the physics and evolution of AGNs and SFGs. The LoTSS is producing high-fidelity images at a central frequency of 144 MHz with a resolution of $6''$ and with decl.-dependent sensitivity, typically around $100 \mu\text{Jy beam}^{-1}$. This is a factor of 10 more sensitive than previous high-resolution sky surveys, e.g., VLA's FIRST. LoTSS will ultimately detect over 10 million radio sources with a significant fraction of star-forming galaxies. A large fraction of LoTSS sources will have optical identifications and photometric redshifts will be available (Duncan et al. 2019; Williams et al. 2019). Furthermore, the WHT Enhanced Area Velocity Explorer multi-object spectrograph on the William Herschel Telescope will observe optical (370–970 nm) spectra of millions of LOFAR radio sources and provide precise redshift information (Smith et al. 2016).

The LoTSS survey, homogeneously covering the whole northern sky complete down to the sub mJy limit will overcome statistical limitations due to shot noise. The large galaxy number density and large sky coverage will substantially reduce cosmic variance in cosmological analysis. The radio galaxies, tracing the background dark matter, will constrain the shape of power spectrum, i.e., the early universe physics, dark



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matter, baryon, density of neutrinos, the inflation power spectrum, and the degree of non-Gaussianity in density fluctuations. The upcoming LoTSS catalogs, covering a large sky area, will help us to explore further regarding large-scale anomalies (de Oliveira-Costa et al. 2004; Ralston & Jain 2004; Schwarz et al. 2004; Tiwari & Aluri 2019) and the current puzzling dipole signal observed with radio catalogs (Blake & Wall 2002; Singal 2011; Gibelyou & Huterer 2012; Rubart & Schwarz 2013; Tiwari & Jain 2015; Tiwari et al. 2015; Tiwari & Nusser 2016; Colin et al. 2017; Siewert et al. 2021). Furthermore, LoTSS will significantly improve on present low-frequency radio catalogs, e.g., TIFR GMRT Sky Survey (TGSS; Intema et al. 2017) and GaLactic and Extragalactic All-sky MWA (GLEAM; Hurley-Walker et al. 2017), and analyses based on these surveys (Dolfi et al. 2019; Rana & Bagla 2019; Tiwari 2019; Tiwari et al. 2019; Choudhuri et al. 2020). Unfortunately, the link between the galaxy and total matter power spectra depends on some unknowns from astrophysics, such as the galaxy bias factor, which depends on galaxy type and is quite different for radio AGNs and star-forming galaxies. The LoTSS population is a mixture of AGNs and star-forming galaxies, and therefore understanding galaxy bias, relative number densities, and luminosity evolution is nontrivial. The purpose of this work is to present a detailed cosmological analysis of LoTSS galaxies and study the effect of survey footprint, shot noise, and other systematics. We have produced galaxy mocks for the survey and have customized and calibrated the data pipeline for galaxy clustering statistics recovery.

We introduce the catalog for LoTSS DR1, as well as its completeness and data mask, in Sections 2 and 2.1. In Section 3, we present mock generation details and covariance matrix estimation. We briefly present the theoretical formulation of the galaxy angular power spectrum and its connection to the dark matter power spectrum in Section 4. In Sections 5 and 6, we present our measured angular power spectrum and two-point correlation statistics, respectively. We present an estimate for the bias for the LoTSS population in Section 7, and summarize our results in Section 8. We conclude with a discussion in Section 9.

2. LoTSS DR1 Catalog

The LOFAR Two-meter Sky Survey (LoTSS; Shimwell et al. 2017) is ongoing and plans to scan the entire northern sky at 120–168 MHz. Shimwell et al. (2019) have prepared the first full-quality public data release (LoTSS-DR1) catalog from 63 LoTSS data sets (2% of the total survey) in the region of the HETDEX Spring Field that were observed between 2014 May 23 and 2015 October 15. The LoTSS DR1 has been prepared using a fully automated direction-dependent calibration and imaging pipeline discussed in Shimwell et al. (2017). The catalog covers 424 square degrees and contains a total of 325,694 sources with peak flux density at least five times the local rms noise, and thus a source density of about 770 sources per square degree. The distribution of radio galaxies in the LoTSS-DR1 catalog is shown in Figure 1. The resolution of the survey images is $6''$ and the positional accuracy of the sources in the catalog is within $0''.2$. The median sensitivity is $71 \mu\text{Jy beam}^{-1}$ at 144 MHz. Williams et al. (2019) remove artifacts, correct wrong groupings of Gaussian components, and prepare a value-added catalog; the catalog then contains 318,520 sources, of which 231,716 have optical/near-IR

identifications in Pan-STARRS⁶/WISE.⁷ Not all of these have photo- z detection, as Pan-STARRS is not as complete as WISE and we only have photo- z measurements for about 50% of all radio sources.

2.1. Completeness and Data Mask

Shimwell et al. (2019) estimate the completeness of LoTSS-DR1 and report the catalog to be 65% complete at 0.23 mJy, 90% at 0.45 mJy, and 95% at 0.58 mJy. The catalog is more than 99% complete at 1 mJy. The LoTSS DR1 catalog was generated by combining 58 individual LOFAR pointings on the sky, and due to poor ionospheric conditions and the presence of bright sources, the imaging pipeline occasionally produces substandard images. As a result, some pointings show inhomogeneous point-source distribution—for example, five pointings in particular show significant incompleteness (Siewert et al. 2020). After properly defining the survey region and masking five incomplete pointings, a reasonable survey mask (mask- z) for reliable cosmology is shown in Figure 2. Mask- z is an upgrade to mask- d (Siewert et al. 2020); it further rejects regions where information from Pan-STARRS is missing. In addition to mask- z , we further consider masking cells with a local noise above the median noise (mask- $z1$) and two times the median noise (mask- $z2$).⁸ We impose the completeness flux cut (i.e., 1 mJy) and also remove the ultra-bright sources with flux equal to or greater than 10 Jy; the catalog thus contains 121,730 sources. After employing mask- z shown in Figure 2, 107,926 sources remain.

3. Mock Catalogs and Covariance Matrix Estimation

Mock catalogs are essential for assessing the analysis pipeline, data systematics and errors in the cosmological analysis of large galaxy surveys. In order to estimate the uncertainty in the cosmological signal recovered from LoTSS data, we generate 1000 LoTSS mocks of the large-scale structure by employing the log-normal (and Gaussian) density field simulator code FLASK⁹ (Xavier et al. 2016). To emulate the LoTSS DR1 catalog, we generate multiple log-normal density fields tomographically in 35 redshift slices (0 to redshift 3.5), each with a width $\Delta z = 0.1$. All statistical properties (i.e., auto- and cross-correlations) are determined by the input angular power spectrum. In addition, effects including redshift-space distortions and lensing are also included in the simulation via the input spectra provided to the FLASK pipeline.

The input theoretical angular power spectra, C_ℓ , in different redshift bins are generated using CAMB (Challinor & Lewis 2011), which includes the effects on the observed number counts of redshift-space distortions, nonlinear power spectrum corrections, and lensing. The redshift distribution profile $N(z)$ is provided along with C_ℓ to FLASK in order to generate mock number count maps. Computation of theoretical C_ℓ requires a galaxy bias, $b(z)$. From our Markov Chain Monte Carlo (MCMC) fits, we iteratively provide the best-fitted bias and generate mock LoTSS DR1 catalogs. We assume a Λ CDM model and use Planck Collaboration et al. (2020) cosmological parameters as the fiducial cosmology throughout this work; in

⁶ <https://www.ifa.hawaii.edu/research/Pan-STARRS.shtml/>

⁷ https://www.nasa.gov/mission_pages/WISE/main/index.html

⁸ We thank Thilo Siewert for preparing and making mask- z , $z1$, and $z2$ available to us.

⁹ <http://www.astro.iag.usp.br/~flask/>

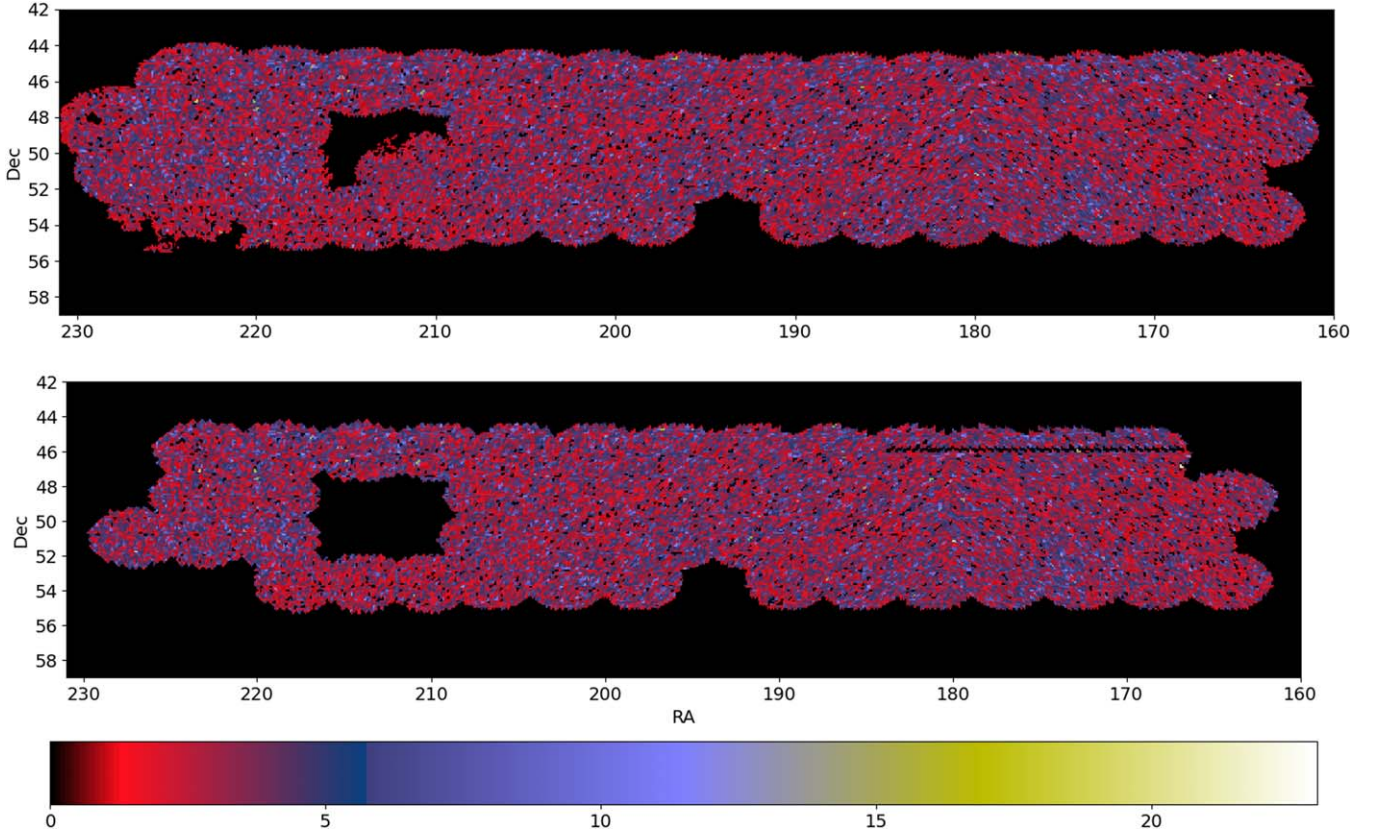


Figure 1. The distribution of radio galaxies in the LoTSS-DR1 catalog with integrated flux density between 1 and 10,000 mJy. Plotted are the number counts in pixels in Cartesian projection at HEALPix (<http://healpix.sourceforge.net>) resolution $N_{\text{side}} = 512$ before (top) and after (bottom) applying the mask. After applying the mask and 1–10,000 mJy flux cut, there are 107,926 sources remaining.

particular, we set the base cosmology parameters from the Planck 2018 baseline, i.e., $\text{TT,TE,EE+lowE+lensing}$ $\omega_b h^2 = 0.02237$, $\omega_c h^2 = 0.1200$, $100\theta_{\text{MC}} = 1.04092$, $\tau = 0.0544$, $\ln(10^{10} A_s) = 3.044$, and $n_s = 0.9649$.

We apply the LoTSS DR1 survey mask to precisely account for survey geometry, then from the mock catalogs we calculate the angular power spectrum up to $l \approx 1500$. Although LoTSS DR1 is a relatively high number density radio catalog, it is hardly over 1% of the sky, and so, given the small sky coverage, C_ℓ is noisy. Therefore, we choose $\Delta\ell = 16$ and recover C_ℓ in bands by collecting 16 multipoles per bin. Finally, with mock C_ℓ , we compute a covariance matrix to determine the uncertainty in LoTSS DR1 recovered galaxy power. A plot of the correlation matrix of the binned angular power spectrum is shown in Figure 3.

4. Galaxy Angular Power Spectrum

The galaxies are biased tracers of the underlying dark matter density and thus the underlying cosmological model. The theoretical relationship between the total matter density perturbations and galaxy density can be easily written in terms of galaxy clustering measures. The statistical measure of clustering of galaxies can be conventionally expressed in terms of the angular power spectrum, C_ℓ . Following the ΛCDM scenario, the theoretical formulation of C_ℓ is briefly described as follows. Assume a uniform galaxy survey catalog with an area \mathcal{A} and total number of galaxies \mathcal{N} . The mean number density, \bar{N} , per steradian is thus simply \mathcal{N}/\mathcal{A} . Subsequently, $\mathcal{N}(\hat{r}) = \bar{N}(1 + \Delta(\hat{r}))$ is the projected number density per

steradian in the direction \hat{r} . Here, $\Delta(\hat{r})$ is the projected number density contrast and is theoretically connected to the underlying matter density contrast, $\delta_m(\mathbf{r}, z(r))$. Note that \mathbf{r} stands for comoving distance r in direction \hat{r} , and $z(r)$ is the redshift corresponding to comoving distance r . The galaxy density contrast, $\delta_g(\mathbf{r}, z(r))$, in the direction \hat{r} and at redshift z , in terms of matter density contrast $\delta_m(\mathbf{r}, z(r))$,

$$\delta_g(\mathbf{r}, z(r)) = \delta_m(\mathbf{r}, z = 0)D(z)b(z), \quad (1)$$

where $b(z)$ is galaxy bias and $D(z)$ is the linear growth factor. Following these, we can formulate the theoretical $\Delta(\hat{r})$,

$$\begin{aligned} \Delta(\hat{r}) &= \int_0^\infty \delta_g(\mathbf{r}, z(r))p(r)dr \\ &= \int_0^\infty \delta_m(\mathbf{r}, z = 0)D(z)b(z)p(r)dr, \end{aligned} \quad (2)$$

where $p(r)dr$, the radial distribution function, is the probability of observing a galaxy between comoving distance r and $(r + dr)$. Note its connection to redshift distribution profile $N(z)$, $p(r) \propto N(z)dz/dr$. The $\Delta(\hat{r})$ may also have some tiny additional contributions from lensing, redshift distortions, and physical distance fluctuations, as well as from variation of radio source luminosities and spectral indices (Chen & Schwarz 2015). However, these effects are expected to be limited to a few percent on the largest scales (Dolfini et al. 2019). Next, we expand $\Delta(\hat{r})$ in terms of spherical harmonics

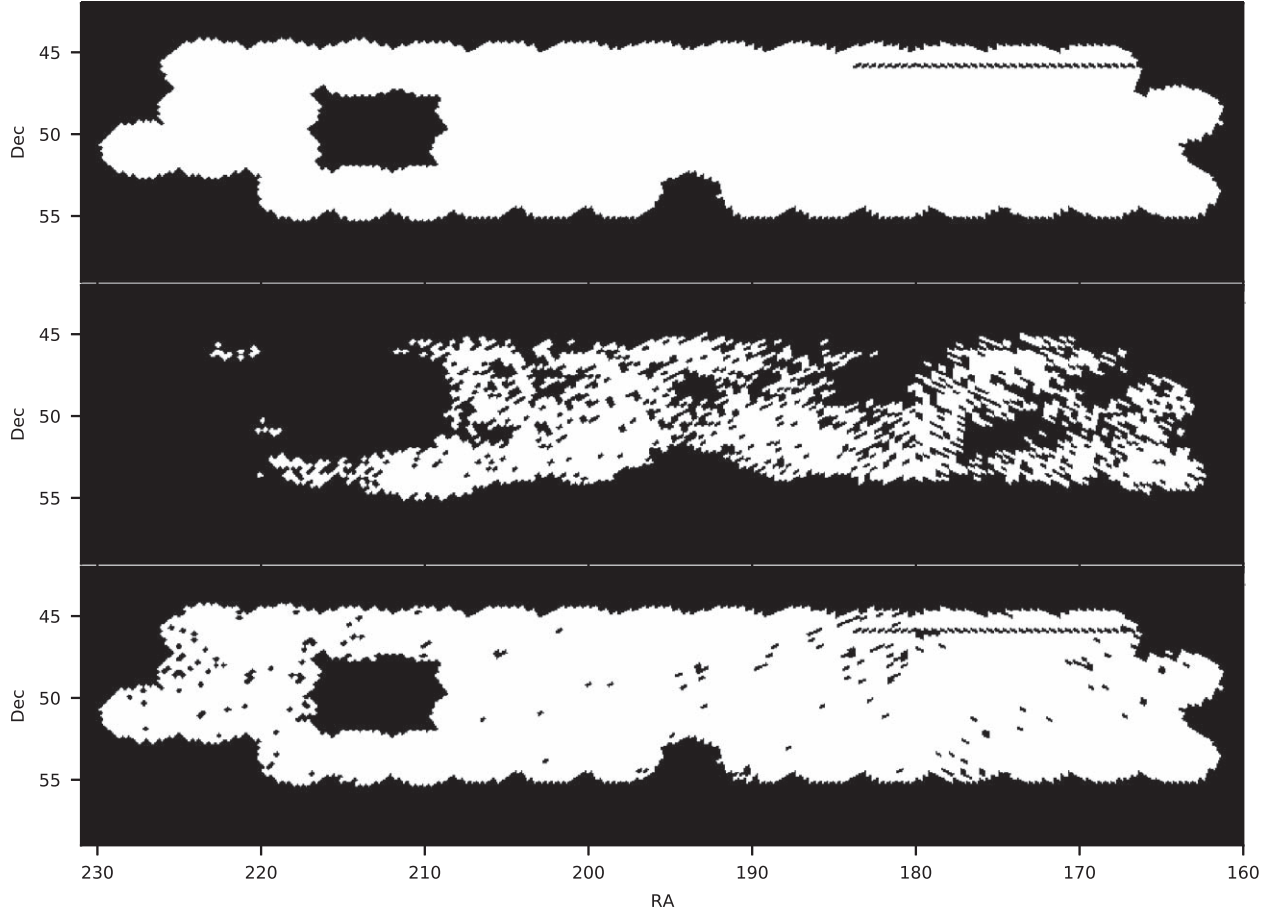


Figure 2. LoTSS DR1 mask-z, z1, and z2. The mask-z (top) rejects five undersampled pointings and includes all 53 good pointings of the LOFAR HETDEX field. It also rejects regions of equal decl., where information from Pan-STARRS is missing. The mask has been prepared considering the survey geometry, consistency of source counts, and sampling uniformity (Siewert et al. 2020). In the middle and bottom panels, we show mask-z1 and mask-z2, which are obtained by further masking cells with a local noise above the median noise and two times the median noise, respectively.

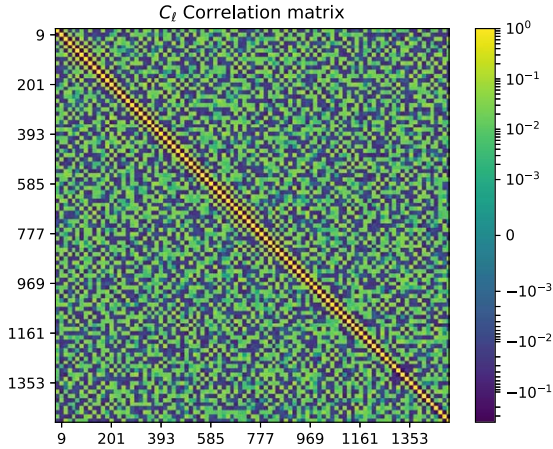


Figure 3. Angular power spectrum correlation matrix constructed from 1000 mocks. The galaxy mocks are generated using the FLASK log-normal density simulator. Each galaxy mock follows the same sky coverage and contains approximately the same number of sources as LoTSS DR1 after mask and flux cuts.

to obtain C_ℓ ,

$$\Delta(\hat{\mathbf{r}}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{\mathbf{r}}). \quad (3)$$

We use the orthonormal property of spherical harmonics and write $a_{\ell m}$ as

$$\begin{aligned} a_{\ell m} &= \int d\Omega \Delta(\hat{\mathbf{r}}) Y_{\ell m}^*(\hat{\mathbf{r}}) \\ &= \int d\Omega Y_{\ell m}^*(\hat{\mathbf{r}}) \int_0^\infty dr \delta_m(\mathbf{r}, z=0) D(z) b(z) p(r). \end{aligned} \quad (4)$$

We can Fourier transform the matter density field $\delta_m(\mathbf{r}, z=0)$ in terms of the \mathbf{k} -space density field $\delta_{\mathbf{k}}$,

$$\delta_m(\mathbf{r}, z=0) = \frac{1}{(2\pi)^3} \int d^3k \delta_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}}, \quad (5)$$

and substitute

$$e^{i\mathbf{k}\cdot\mathbf{r}} = 4\pi \sum_{\ell, m} i^\ell j_\ell(kr) Y_{\ell m}(\hat{\mathbf{r}}) Y_{\ell m}^*(\hat{\mathbf{k}}), \quad (6)$$

where j_ℓ is the spherical Bessel function of the first kind for integer ℓ . Next, we can write $a_{\ell m}$ as

$$a_{\ell m} = \frac{i^\ell}{2\pi^2} \int dr D(z) b(z) p(r) \int d^3k \delta_{\mathbf{k}} j_\ell(kr) Y_{\ell m}^*(\hat{\mathbf{k}}). \quad (7)$$

Subsequently, we find the expression for the angular power spectrum, C_ℓ , as

$$\begin{aligned} C_\ell &= \langle |a_{\ell m}|^2 \rangle \\ &= \frac{2}{\pi} \int dk k^2 P(k) \left| \int_0^\infty dr D(z) b(z) p(r) j_\ell(kr) \right|^2 \\ &= \frac{2}{\pi} \int dk k^2 P(k) W^2(k), \end{aligned} \quad (8)$$

where $P(k)$ is the matter power spectrum, and $W(k) = \int_0^\infty dr D(z) b(z) p(r) j_\ell(kr)$ is the k -space window function. The galaxies are discrete point sources and thus the measured angular power spectrum of a galaxy catalog incorporates the Poissonian shot-noise contribution σ_0 , equal to $\frac{1}{N}$. In addition, a radio catalog in general contains multiple entries for some sources, causing a constant offset to C_ℓ . This offset can be approximated in terms of shot-noise σ_0 . Therefore, the corrected power spectrum, $C_\ell^{\text{corrected}}$, corresponding to the theoretical C_ℓ given in Equation (8) is $C_\ell^{\text{measured}} - \sigma_0 \times \alpha$, where α is 1 (shot noise) + an offset factor due to the multicomponent correction. We have discussed this further in Section 5. The uncertainty in C_ℓ determination due to cosmic variance, sky coverage, and shot noise is

$$\Delta C_\ell = \left(\frac{2}{(2\ell + 1)f_{\text{sky}}} \right)^{1/2} C_\ell^{\text{measured}}, \quad (9)$$

where f_{sky} is the fraction of sky covered by survey.

5. C_ℓ Measurements

We make use of a pseudo- C_ℓ recovery algorithm by Alonso et al. (2019) to achieve an efficient and reasonably accurate estimate of the angular power spectrum of the LoTSS DR1 catalog. The algorithm Python module is publicly available as NaMaster,¹⁰ and the mathematical background of the estimator, its features, and software implementation are described in Alonso et al. (2019). We test the performance of the pseudo- C_ℓ algorithm for the LoTSS DR1 mask using a test map, which resembles the LoTSS DR1 density fluctuations as follows. We obtain a reasonable recovery of C_ℓ by using an un-apodized mask shown in Figure 2. To emulate the LoTSS DR1 catalog, we first generate $a_{\ell m}$ using LoTSS DR1 model C_ℓ ,¹¹ subsequently, we obtain a density contrast map from $a_{\ell m}$ by making use of spherical harmonics (Equation (3)). Next, we apply the LoTSS DR1 mask to this test map and recover C_ℓ (pseudo- C_ℓ). The results thus obtained with different configurations and settings are shown in Figure 4, and above $\ell > 100$ a reasonable recovery is demonstrated.

With the above demonstrated settings, we run the pseudo- C_ℓ recovery algorithm and obtain the LoTSS DR1 catalog angular power spectrum. The angular spectrum thus obtained is shown in Figure 5. The recovered power spectrum for galaxies with integrated radio flux above survey completeness, i.e., $S > 1$ mJy, approximately agrees with theoretical results obtained following Λ CDM and considering NVSS (Condon et al. 1998) radio galaxy biasing and radial distribution (Nusser & Tiwari 2015). Without considering survey completeness (i.e.,

the flux cut), the recovered power spectrum is significantly high at low ℓ , i.e., at large scales, presumably attributing the effect of survey incompleteness. At low fluxes, when the survey is not complete, we may have large-scale flux systematics and uneven source counts, and thus may observe excess large-scale clustering. Considering a slightly higher flux cut, i.e., $S > 2$ mJy, we obtain similar results as for the $S > 1$ mJy flux cut C_ℓ but with slightly more fluctuations. The results are shown in Figure 6. We next test the robustness of C_ℓ recovery with different masks discussed in Section 2.1. With mask-z, z1, and z2 we recover almost the same C_ℓ , although the fluctuations for mask-z1 are slightly higher due to low sky coverage (i.e., high cosmic variance). The results are shown in Figure 7.

Depending on the flux density threshold considered, the LoTSS DR1 catalog contains a significant number of multicomponent sources (Siewert et al. 2020). The multicomponent sources cause a constant offset to C_ℓ . Siewert et al. (2020) estimate the clustering parameter n_c for LoTSS DR1, which is defined to be the ratio of the variance of counts in cells over the mean of the counts in cells. In addition, they notice that the counts-in-cells statistic follows a compound Poisson distribution, assuming that the parameter $n_c = 1 + \gamma$, where γ is the mean number of components for a source. Given this, for any source in the catalog, the probability of having one or two components can be given as $\gamma \exp(-\gamma)$ or $\frac{\gamma^2}{2!} \exp(-\gamma)$, respectively. This implies that the ratio between the number of two-component and one-component sources equals $\gamma/2$, and thus the fraction of radio sources split into double sources, i.e., e as defined in Blake et al. (2004), equals $\frac{\gamma/2}{1+\gamma/2}$. We further note from Siewert et al. (2020) that $\gamma = 0.44$ for 1 mJy flux density threshold, and estimate the multicomponent offset in terms of shot noise σ_0 to equal $\gamma\sigma_0/(1+\gamma) \approx 0.3\sigma_0$ for C_ℓ following Blake et al. (2004). The estimate of course makes several assumptions, as discussed above, and is not guaranteed to be exact. We thus consider the multicomponent offset to C_ℓ as an independent parameter in our fitting procedure. Indeed, we recover a multicomponent offset close to $0.3\sigma_0$, as expected.

6. Two-point Correlation Statistics

For comparison with theory and earlier results in Siewert et al. (2020), we calculate the two-point correlation function (2PCF) for the LoTSS DR1 catalog. We obtain approximately the same 2PCF results with our catalog and mask, as in Siewert et al. (2020) with 1 mJy flux cut and their mask-d. The theoretical curve at angular scales less than 0.1° is slightly lower than the measured values, presumably due to the multiple-component nature of some radio sources in the catalog (Blake et al. 2004). We have followed Siewert et al. (2020) and applied the same Python package TreeCorr¹² (Version 4.0) (Jarvis et al. 2004) and the same parameter settings for 2PCF estimation. Following Siewert et al. (2020), we next divide the LoTSS DR1 catalog into three patches: “left,” “center,” and “right.” The R.A. ranges for these patches are (161, 184), (184, 208), and (208, 230) degrees, respectively. The calculated angular correlation functions are shown in Figure 8. We find that the center area fits the theory best. The slight mismatch between the “left,” “center,” and “right” patch results is an indication of large-scale density fluctuation systematics present in the data. This should be considered as a motivation for refining the LoTSS DR1 pipeline (Shimwell et al. 2019).

¹⁰ <https://namaster.readthedocs.io/en/latest/index.html>

¹¹ C_ℓ is the expected variance of the $a_{\ell m}$ at ℓ .

¹² <http://github.com/rmjarvis/TreeCorr/>

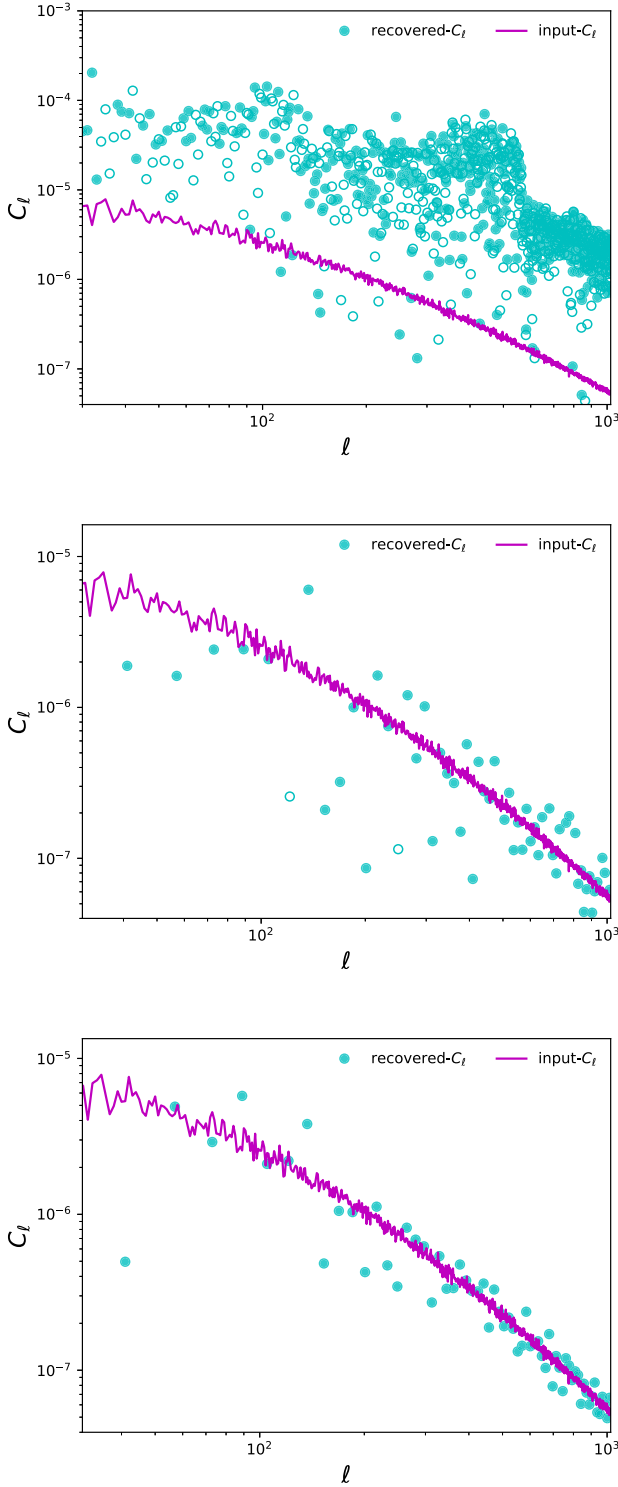


Figure 4. Angular power spectrum recovery performance with LoTSS DR1 mask (0.1% of the sky). Empty circles show a negative value. The recovered power spectrum is significantly fluctuating as a function of multipole (top figure). Next, we apodized the mask on the scale of one degree and plot C_ℓ in bands by collecting 16 multipoles per bin; we show the results for this in the middle panel. The best recovery above $\ell > 100$ is obtained if we collect 16 multipoles per bin and do not apodize the mask (bottom panel).

Large-scale systematics in real space correspond to low multipoles of angular power spectrum C_ℓ . Therefore, to examine the effect more closely, we calculate the angular

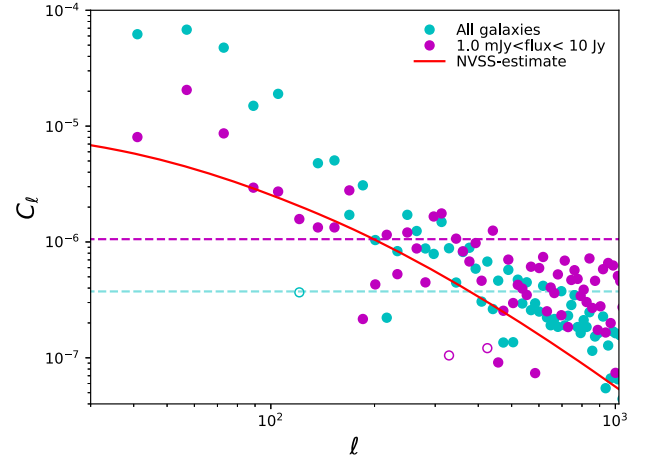


Figure 5. LoTSS DR1 angular power spectrum recovery. We obtain an approximate theoretical power spectrum “NVSS-estimate” for the LoTSS DR1 catalog following Equation (8) and assuming NVSS galaxy bias and redshift distribution (Nusser & Tiwari 2015). The power spectrum obtained without a flux cut is significantly high at large scales (low ℓ). The dashed lines show the shot noise. The LoTSS DR1 catalog is only expected to be complete above 1 mJy.

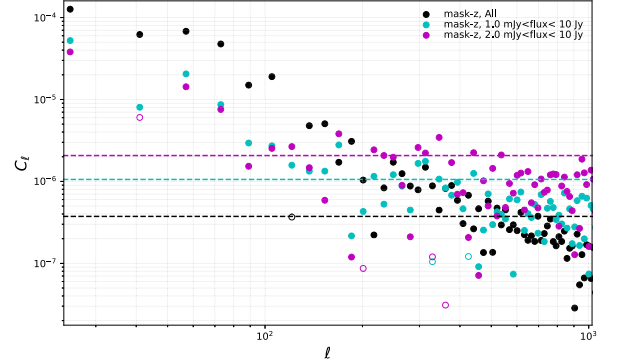


Figure 6. LoTSS DR1 angular power spectrum recovery with different flux cuts. The empty circles show the negative value. The dashed lines show the shot noise.

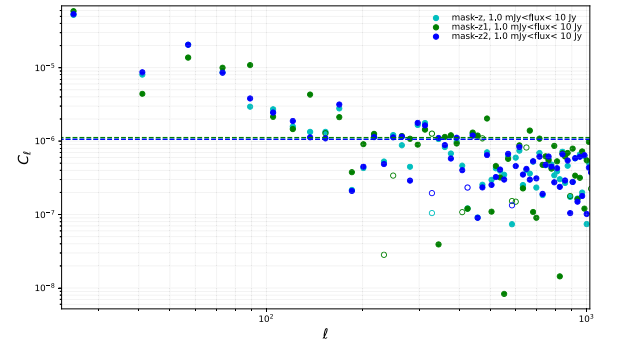


Figure 7. LoTSS DR1 angular power spectrum recovery with mask-z, z1, and z2. These masks are described in Section 2.1 and shown in Figure 2. Other details same as Figure 6.

power spectrum for these three R.A. patches. We show the results in Figure 9 and note that C_ℓ recoveries from different R.A. patches are almost the same above $\ell > 10$. We recognize that $\ell = 10$ corresponds to a scale $\approx 20^\circ$, which is approximately the R.A. patch width considered. We note that we only

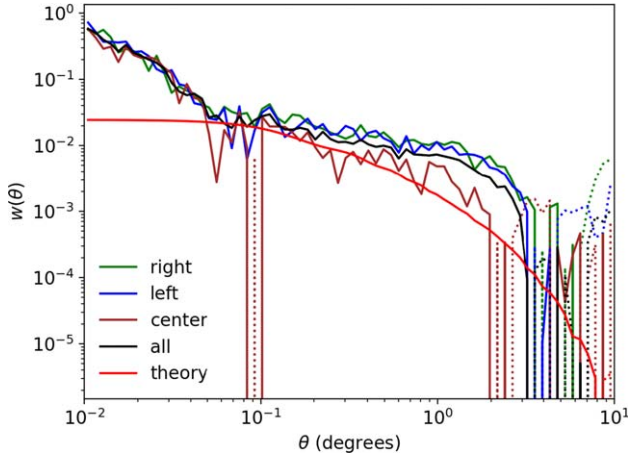


Figure 8. The two-point correlation function of LoTSS DR1 catalogs in different patches of the sky. The green curve denotes the “right” patch angular correlation function, blue the “left” area, and orange the “center” area. The whole area angular correlation function of LoTSS DR1 is shown by the black line, while the red curve denotes the theoretical prediction following Λ CDM with NVSS bias and redshift distribution. All dotted lines represent negative values.

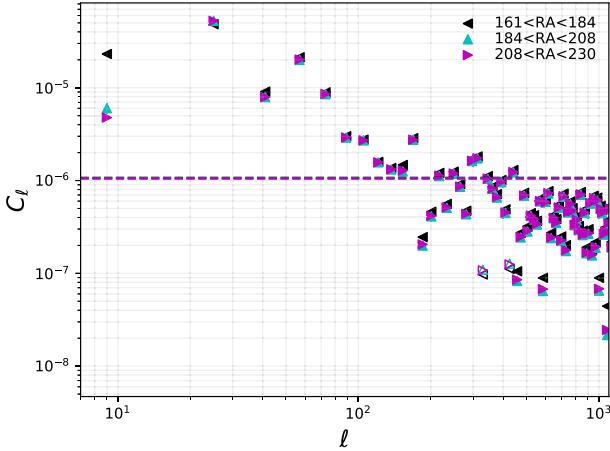


Figure 9. Angular power spectrum recovery of LoTSS DR1 catalogs in different patches of the sky. The recovered power spectrum from different R.A. patches is approximately the same above $\ell > 10$. The dashed lines show the shot noise.

consider data above $\ell = 100$, and for this ℓ range we do not see any large-scale R.A. dependence.

7. Galaxy Bias and $N(z)$ Estimate

The bias $b(z)$ for each galaxy population is different and depends on the dark matter halo mass hosting the galaxy type (Mo & White 1996). The angular power spectrum of LoTSS DR1 in Figure 4 is apparently a close match with Λ CDM estimates considering NVSS galaxies bias and redshift distribution, $N(z)$. However, in reality, the bias and $N(z)$ for LoTSS DR1 galaxies could be significantly different in comparison with values obtained for NVSS galaxies with flux above 10 mJy (NVSS completeness) at 1.4 GHz (Nusser & Tiwari 2015; Tiwari & Nusser 2016). The NVSS is dominantly radio-loud AGNs, Fanaroff–Riley type I (FR I) and FR II (Fanaroff & Riley 1974), whereas the LoTSS DR1 above 1 mJy contains star-forming galaxies along with radio-loud AGNs, FR I, and FR II type (Wilman et al. 2008; Calistro Rivera et al. 2017). The LoTSS DR1 population is mixed and differs from

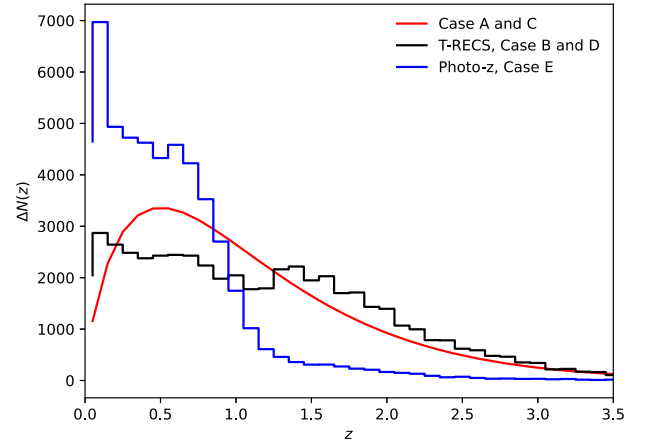


Figure 10. Number of radio sources per redshift bin $\Delta z = 0.1$ of available z for flux density thresholds $S > 1$ mJy. The redshift distributions considered for fitting Λ CDM cosmology are also shown.

NVSS, and thus the bias $b(z)$ and $N(z)$ could be significantly different. For $N(z)$, around half of the LoTSS DR1 sources have optical/near-IR identification in Pan-STARRS/WISE and their photometric redshifts are available. The number count of these photometric galaxies in redshift bins ($\Delta z = 0.1$) is shown in Figure 10. Note that the $N(z)$ in Figure 10 only represents part of the LoTSS DR1 population; the true $N(z)$ of the full LoTSS DR1 catalog could be significantly different. Nevertheless, the available photoredsifts can be assumed to loosely represent the LoTSS DR1 galaxies’ redshift distribution. Furthermore, there are a few more reasonable $N(z)$ templates we may assume for LoTSS DR1, namely the $N(z)$ obtained for NVSS in Nusser & Tiwari (2015) and the $N(z)$ presented in the Tiered Radio Extragalactic Continuum Simulation (T-RECS; Bonaldi et al. 2019). Assuming a Λ CDM cosmology, we seek to constrain the galaxy bias $b(z)$ and radial distribution $N(z)$ of LoTSS DR1 galaxies by fitting the measured angular power spectrum in Figure 5 and the observed $N(z)$ photo- z histogram in Figure 10. For a given bias, $b(z)$, and $N(z)$, we calculate the theoretical (model) angular power spectrum C_ℓ^{th} . Next, we obtain the likelihood, i.e., the probability to have observed data $C_\ell^{\text{corrected}}$ given the model, $\mathcal{P}(C_\ell^{\text{corrected}}|b(z), N(z)) \propto \mathcal{L}(C_\ell^{\text{corrected}}) \propto \exp(-\frac{1}{2}(v_b C_\ell^{\text{corrected}} - v_b C_\ell^{\text{th}})^T v_b \Sigma^{-1} (v_b C_\ell^{\text{corrected}} - v_b C_\ell^{\text{th}}))$. Here, $v_b \Sigma$ is the covariance matrix computed using mocks discussed in Section 3. We make use of Bayes’ probability theorem,¹³ and write the model probability given the data, $\mathcal{P}(b(z), N(z)|C_\ell^{\text{corrected}}) = \mathcal{P}(C_\ell^{\text{corrected}}|b(z), N(z)) \times \mathcal{P}(b(z))$. Here, $\mathcal{P}(b(z))$ is the prior probability of bias $b(z)$. We assume different $N(z)$ templates and run an MCMC sampler to maximize model parameter probability, $\mathcal{P}(b(z), N(z)|C_\ell^{\text{corrected}})$, i.e., to obtain the best fit to Λ CDM cosmology and deduce the effective bias $b(z)$, given the $N(z)$ values. We fit one more model parameter, α , to account for the multicomponent offset to $C_\ell^{\text{corrected}}$ as discussed in Sections 4 and 5. We conveniently use Cobaya (Torrado & Lewis 2021) to perform Bayesian analysis and CosmoMC (Lewis & Bridle 2002; Lewis 2013) for MCMC sampling. With mock catalogs, we verify the above MCMC pipeline and recover the input bias within 1σ uncertainty for the $\ell = 200$ to 1000 range. The pipeline performance is demonstrated in Figure 11. Above

¹³ $P(A|B)P(B) = P(B|A)P(A)$.

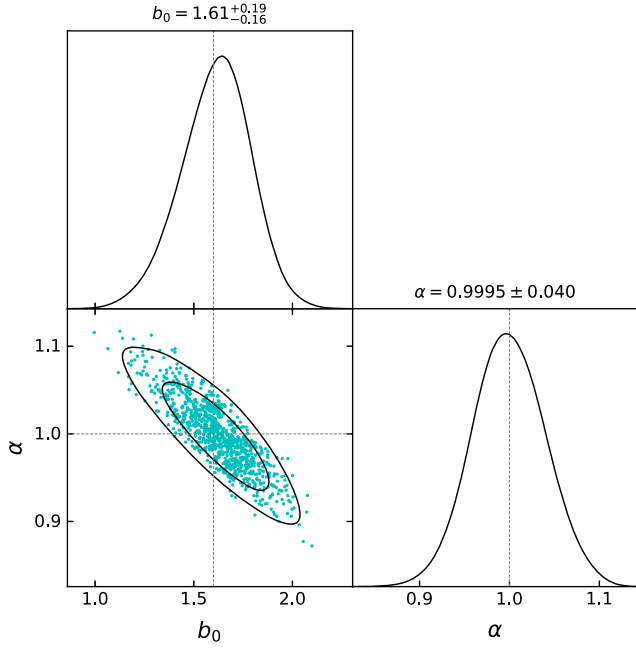


Figure 11. The pipeline recovers input bias parameters within 1σ ; dashed lines are input values. The figure in particular demonstrates the quadratic bias $b(z) = b_0 + b_1z + b_2z^2$ recovery. For our fitting range, i.e., $\ell = 200\text{--}1000$, we do not cover enough power spectrum shape to constrain b_1 and b_2 , and the posterior probabilities for these are almost the same as the prior information we assume.

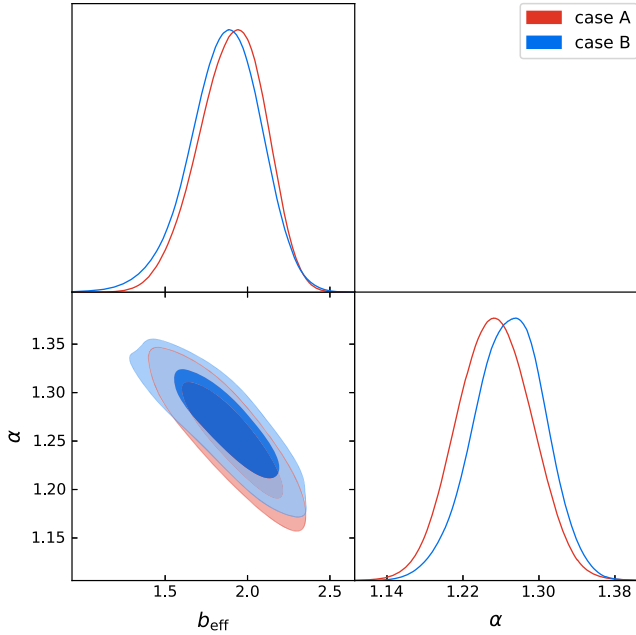


Figure 12. Cases A and B: Posterior distributions of bias and shot-noise factor α . Here, $(\alpha - 1) \times \sigma_0$ is the multicomponent offset.

$\ell = 1000$, the recovery from mocks is noisy, and therefore we restrict our fits to $\ell \leq 1000$ for all cases discussed below. We would like to emphasize that the measured C_ℓ are projected quantities and it is hard to extract tomographic information from them. The recovered bias is strongly dependent on the $N(z)$ assumption; the best we can do for now is to assume a variety of $N(z)$ models for the LoTSS catalog. We hope to have better $N(z)$ observations in future and thus better constraints on $b(z)$.

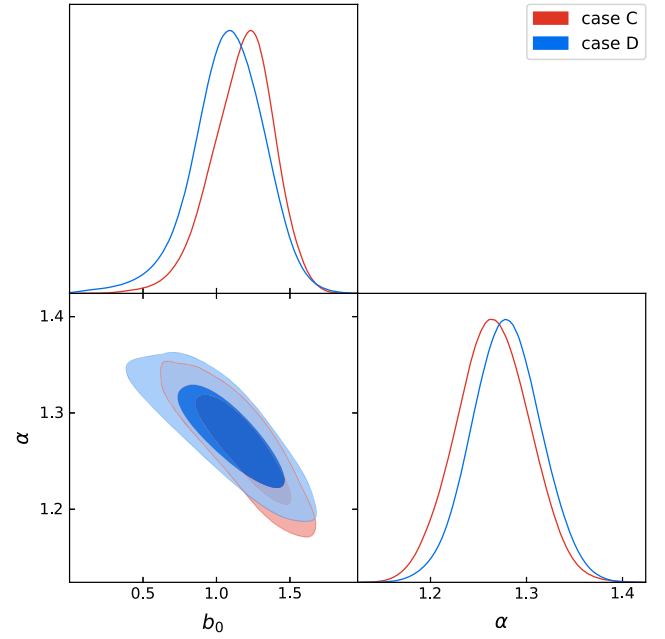


Figure 13. Case C and D: One- and two-dimensional posterior distributions of quadratic-bias-free parameters b_0 and shot-noise factor α . As the data is not sensitive to b_1 and b_2 , we fix these to NVSS observed values. Other details same as in Figure 12.

Case A: Assume $N(z) \sim z^{0.74} \exp\left[-\left(\frac{z}{0.71}\right)^{1.1}\right]$ as given in Nusser & Tiwari (2015); we fit an effective redshift-independent bias and shot-noise factor α . The posterior distributions of bias and shot-noise factor α are shown in Figure 12. We find that a bias $b_{\text{eff}} = 1.90 \pm 0.20$ fits best with the data. The recovered best-fit angular power spectrum is shown in Figure 16. The recovered shot-noise factor is 1.25 ± 0.04 , which is consistent with 1.30 as expected in Section 5.

Case B: We employ the T-RECS (Bonaldi et al. 2019) simulation, and produce an $N(z)$ histogram considering the SFG +AGN population above 1 mJy flux at 150 MHz. The histogram is shown in Figure 10. We assume the T-RECS histogram¹⁴ in Figure 10 represents the LoTSS population and run the MCMC sampler to find the best redshift-independent bias value. The posterior distributions from MCMC sampler are shown in Figure 12. We obtain $b_{\text{eff}} = 1.86 \pm 0.22$ as the best fit to the data. The best-fit angular power spectrum and bias with 1σ uncertainty band are shown in Figures 16 and 15, respectively.

Case C: $N(z)$ is the same as in Case A; we fit a quadratic bias $b(z) = b_0 + b_1z + b_2z^2$. This is to test how well the NVSS quadratic bias obtained in Nusser & Tiwari (2015) fits to LoTSS DR1. We have noticed with mocks that the data is not sensitive to b_1 and b_2 ; we recover the same posterior probability for these as the prior information we assume. We fix these to NVSS observed values, i.e., $b_1 = 0.85$, $b_2 = 0.33$, and run the MCMC sampler over b_0 and noise+multicomponent factor α . The best-fit bias, $b(z)$, and angular power spectrum thus recovered are shown in Figures 15 and 16, respectively. The MCMC sample and posterior distributions are shown in Figure 13. We do see a better fit with quadratic bias; the fit parameters AIC and BIC are slightly improved. All parameters, their prior information, and the fit parameters are given in Table 1.

¹⁴ A cubic spline interpolation is used to make $N(z)$ smoother.

Table 1
 MCMC Sampler Results

		Case A	Case B	Case C	Case D	Case E
b_0	prior			[0.0, 5.0]	[0.0, 5.0]	
	fit value	1.18 ± 0.22	1.08 ± 0.25	...
b_1	fix			0.85	0.85	
b_2	fix			0.33	0.33	
b_{eff}	prior	[0.5, 5.0]	[0.1, 5.0]			[0.1, 5.0]
	fit value	1.90 ± 0.20	1.86 ± 0.22	1.05 ± 0.10
α	prior	[0.01, 3.0]	[0.01, 3.0]	[0.01, 3.0]	[0.01, 3.0]	[0.01, 3.0]
	fit value	1.25 ± 0.04	1.27 ± 0.04	1.26 ± 0.04	1.28 ± 0.04	1.22 ± 0.04
χ^2/dof		1.18	1.36	1.18	1.36	1.39
AIC		60.65	69.28	60.49	69.06	70.66
BIC		64.47	73.11	64.32	72.88	74.48

Notes. The prior and maximum likelihood value of the model free parameters and uncertainties. The flat prior range is shown in square brackets. The reduced χ^2 , i.e.,

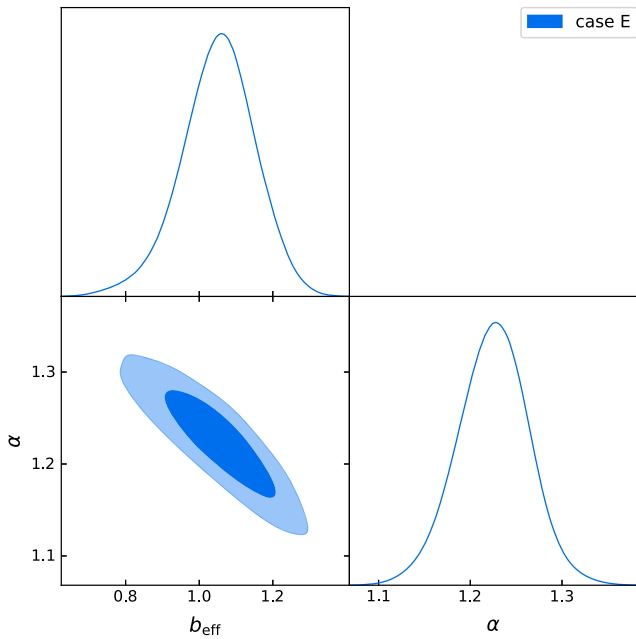


Figure 14. Case E: Posterior distributions of bias and shot-noise factor α . Other details are the same as in Figure 12.

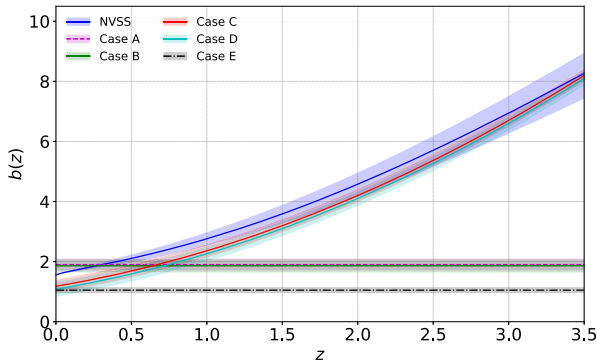


Figure 15. The bias, $b(z)$, recovered with MCMC Bayesian fitting for different cases. The best bias for NVSS radio AGNs obtained in Nusser & Tiwari (2015) is also shown.

Case D: We choose $N(z)$ to be the same as in Case B, and fit a quadratic bias $b(z) = b_0 + b_1z + b_2z^2$. Other details are the same as in Case C. The recovered best bias, $b(z)$, and angular

χ^2 per degree of freedom, AIC, and BIC are also listed for model fit comparison.

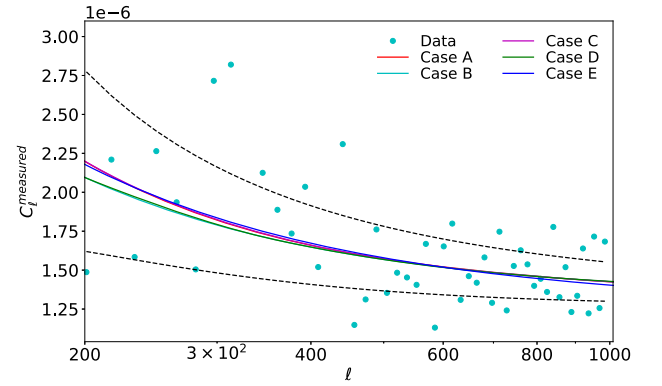


Figure 16. The best-fit Λ CDM C_ℓ^{measured} obtained using MCMC sampling with various sets of bias, $b(z)$, and $N(z)$ constraints and assumptions as discussed in Section 7. The estimated observed pseudo- C_ℓ from data are also shown. The black dashed curves show the 1σ limit due to shot noise and cosmic variance following Equation (9) for Case A. For other cases, the 1σ limit curves are approximately the same.

power spectrum are shown in Figures 15 and 16, respectively. The MCMC sample and posterior distributions are shown in Figure 13.

Case E: Finally, we use observed photo- z redshifts, which are available for only about half of the radio sources. This $N(z)$ only represents part of the LoTSS DR1 population, but we include this case to check the sensitivity of the data to $N(z)$. We assume a redshift-independent bias, i.e., $b(z) = b_{\text{eff}}$, and a shot-noise factor α , and then fit with the recovered LoTSS DR1 angular power spectrum. The posterior distributions from MCMC sampler are shown in Figure 14. We obtain a best-fit bias $b_{\text{eff}} = 1.05 \pm 0.10$. The recovered best power spectrum is shown in Figure 16.

8. Summary

We have explored the first data release of the LoTSS survey for cosmological studies. We have generated detailed data mocks and constructed the covariance matrix for angular power spectrum recovery. With our mocks, we were able to customize the settings of the pseudo-angular power recovery algorithm for LoTSS. The LoTSS DR1 contains dominantly star-forming galaxies and a significant fraction of AGNs; therefore, the resultant bias and radio distribution of galaxies cannot be naively assumed to be the same as seen with other radio

surveys with high flux limits, e.g., the NVSS, where the catalog almost entirely consists of AGNs. We have explored the galaxy bias and possible radial distribution profiles for LoTSS galaxies and have presented their fits to data. We find that the data above integrated flux density 1 mJy and with mask-z has a reasonable angular power spectrum. Furthermore, we consider the multicomponent fraction of radio galaxies as an independent fit parameter and recover its value. For convenience, the summary of different cases, the recovered bias with fit quality parameters, i.e., χ^2/dof , and model selection criteria parameters, the Akaike information criterion (AIC), and the Bayesian information criterion (BIC) are listed in Table 1.

We have tried fitting the best bias and $N(z)$ profiles from NVSS (Nusser & Tiwari 2015; Tiwari & Nusser 2016) to LoTSS, and have found that the bias for LoTSS DR1 is slightly lower in comparison with NVSS if we assume NVSS $N(z)$ (Case C). Our bias values for cases A, B, C, and D are similar to Alonso et al. (2021), obtained using CMB lensing and LoTSS DR1 data. Considering redshift information dominantly determined by available photo- z information (Case E), we obtain a lower bias value close to 1; this low value of bias presumably demonstrates that (i) the data are sensitive to the $N(z)$ distribution and (ii) the photo- z distribution for the limited number of sources does not represent a reasonable $N(z)$ for the full LoTSS DR1 population. Considering different $N(z)$ templates and bias $b(z)$ forms, with an MCMC sampler, we find that the best-fit result (best reduced- χ^2 , AIC, and BIC) is found with Case C, i.e., by assuming $N(z)$ given in Nusser & Tiwari (2015) and a redshift-dependent quadratic bias. The reduced χ^2 of all models is slightly higher than 1. This may indicate some residual systematics. We expect LoTSS DR2 to have better control of systematics and thus to provide better constraints on cosmology and galaxy bias. The multicomponent fraction for LoTSS DR1 is not very settled, and we consider this as an independent parameter. We find that a constant offset to C_ℓ equal to ≈ 0.25 times shot noise is needed. The performance of our MCMC sampler has been verified by recovering model input bias, $b(z)$, by fitting the pseudo- C_ℓ mean from mocks.

To compare our results with Siewert et al. (2020), we have also presented measurements for the 2PCF for our catalog. We obtain approximately the same results as in Siewert et al. (2020). The 2PCFs calculated for broad R.A. patches differ; this potentially indicates large-scale systematics in the data. These large-scale anomalies in the real space clustering signal correspond to low multipoles, approximately $\ell < 10$. For our fits, we have considered the power spectrum above multipole $\ell = 200$, and for this ℓ range, we have verified that the recovered power spectrum is approximately the same for all three R.A. patches.

9. Conclusion and Discussion

We have presented detailed mocks and projected clustering statistics of point sources from LoTSS first data release, based on the observation of 424 square degrees of the sky at ~ 150 MHz. We assume different possible redshift templates, and run the MCMC sampler to obtain the best galaxy bias and multicomponent fraction for LoTSS DR1 above 1 mJy. With two-point correlation statistics, we notice somewhat different clustering in large R.A. patches, presumably indicating large-scale flux calibration systematics in data. The angular power





spectrum C_ℓ , the spherical harmonic decomposition of the fluctuations in angular space, efficiently separates the scale-dependent real space number density fluctuations. The catalog contains multiple entries for a significant number of sources, resulting in a constant offset to the measured angular power spectrum. With the multicomponent correction, we see a good match of LoTSS DR1 measured C_ℓ with standard Λ CDM theory. The data above 1 mJy and with mask- z result in consistent C_ℓ s and arguably qualify for cosmological analysis.

The LOFAR observations are at relatively low radio frequency, and the strength of the systematic effects affecting the data is significantly higher than in high-frequency surveys, e.g., VLA NVSS, and thus the calibration and source cataloging are challenging. Even so, the data calibration pipeline by Shimwell et al. (2019) delivers a workable source catalog for cosmology. The LOFAR team will continue to improve on flux calibration and cataloging. Furthermore, we need better estimates for redshifts; the present partial availability of photometric redshifts is not reliably representing the LoTSS population. Ideally, above the LoTSS completeness limit, we need accurate redshifts for all galaxies over a small sky patch to obtain a complete redshift distribution of LoTSS galaxies (Duncan et al. 2021). With improved calibration, large sky coverage, and higher number density, the upcoming catalogs from LOFAR are expected to be a significant improvement over current low-frequency radio catalogs, e.g., the TGSS from GMRT and GLEAM from MWA. The upcoming LoTSS catalogs are an excellent opportunity for cosmological studies in the low radio frequency regime.

We thank Thilo Siewert for helping with the LoTSS DR1 catalog and masks. P.T. acknowledges the support of the RFIS grant (No. 12150410322) by the National Natural Science Foundation of China (NSFC). P.T., R.Z., J.Z., and G.B.Z. are supported by the National Key Basic Research and Development Program of China (No. 2018YFA0404503) and NSFC Grants 11925303 and 11720101004, and a grant of CAS Interdisciplinary Innovation Team. J.Z. is supported by Chinese Scholarship Council (CSC) and Science and Technology Facilities Council (STFC) for visiting University of Portsmouth. Some of the results in this paper have been derived using the HEALPix (Górski et al. 2005) package.

LOFAR data products were provided by the LOFAR Surveys Key Science project (LSKSP; <https://lofar-surveys.org/>) and were derived from observations with the International LOFAR Telescope (ILT). LOFAR (van Haarlem et al. 2013) is the Low Frequency Array designed and constructed by ASTRON. It has observing, data processing, and data storage facilities in several countries, which are owned by various parties (each with their own funding sources), and which are collectively operated by the ILT foundation under a joint scientific policy. The efforts of the LSKSP have benefited from funding from the European Research Council, NOVA, NWO, CNRS-INSU, the SURF Co-operative, the UK Science and Technology Funding Council, and the Jülich Supercomputing Centre.

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