

1 **Stress-Dependent Fluid Dynamics of Shale Gas Reservoirs: A**
2 **Pore Network Modeling Approach**

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11 **Abstract**

12 In this study, the pore network modeling approach is implemented to study the flow behavior
13 of the methane gas in shale reservoirs in the presence of slip flow and gas desorption under
14 different effective stress conditions through the effect on apparent permeability. A single pore
15 and a pore network are constructed using a down-scaled Berea sandstone image. The
16 compressible Stokes is applied with proper boundary conditions to capture slip flow and gas
17 desorption effects. To include the effect of stress, poroelasticity principles are applied to
18 describe the interaction between fluid flow and solid deformation. The model performance is
19 verified against an experimental data taken from the literature. Both models show a negative
20 effect of stress on gas flow when slip flow and desorption mechanisms are activated in the
21 models. In the single pore model, slip has a positive and complementary effect on desorption
22 at a specific stress value when their coupled contributions are summative. However, the results

23 of pore network model show that slip has a destructive effect on desorption contribution
24 specially at low effective stress values. That is, the changes in pore pressure induced by the
25 slip mechanism act to limit the desorption contribution. This is clearly noted with the
26 consideration of stress effects. At high effective stress values, the slip effect is limited due to
27 the increase in pore pressure by compression, hence a closer matching between the
28 permeability enhancement produced by the desorption mechanism alone and the coupled effect
29 of desorption and slip can be noticed. Stress acts to limit the slip flow and also the desorption
30 mechanism as they both are a function of pore pressure, however, as stress increases the
31 contribution of desorption no longer would be reduced by slip flow. This study provides
32 insights into stress role in multiphysics gas flow in shale reservoirs.

33 **Keywords:** Pore network modeling; Shale reservoirs; Gas desorption; Slip flow; Effective
34 stress; Apparent permeability

35 **1. Introduction**

36 Shale gas reservoirs are categorized under unconventional resources, where low porosity,
37 ultra-low permeability, nanoscale pore and pore throat radius, stress dependency of the inherent
38 properties, and complex gas flow mechanisms are listed as the technical issues of these
39 reservoirs [1]. Despite the mentioned complexities, a large amount of natural gas all around the
40 world is produced from shale gas reservoirs, especially in the United States (34% based on the
41 EIA report in 2011) [2].

42 The permeability of the shale gas reservoirs is within the range of ultra-low (several nano-
43 Darcy to several hundred nano-Darcy) and it varies with the effective stress of porous media.
44 This range of permeability is representative of the nanoscale pore and pore throat size, which

45 are close to the mean free path of the molecules. Thus, several flow regimes, such as continuum
46 flow, slip flow, transitional flow, and molecular-free flow, are expected to occur during the
47 different stages of production [3]. The mentioned flow regimes are classified using the
48 Knudsen number, where it describes the frequency collisions of molecule-molecule and
49 molecule-wall based on the theory of the gas molecular dynamics [4, 5], as presented in Table
50 1.

51 Table 1. The flow regimes and their corresponding Knudsen numbers [3]

Flow regime	Knudsen number (K_n)
Continuum flow	$0 - 10^{-3}$
Slip flow	$10^{-3} - 10^{-1}$
Transition flow	$10^{-1} - 10^{+1}$
Free-molecular flow	$10^{+1} - \infty$

52 As mentioned above, different flow regimes occur during the different stages of the
53 production from a shale gas reservoir that can significantly affect its apparent permeability.
54 The expression between gas slip flow phenomena and apparent permeability was introduced
55 by Klinkenberg for the first time [6]. After that, a new line of research on the relation between
56 slip flow and apparent permeability has been flourished, while in recent years it was studied by
57 many researchers. For instance, Beskok and Karniadakis (1999) described different flow
58 regimes by a comprehensive model based on the Knudsen number and the Hagen-Poiseuille
59 equation. The proposed model can cover continuum flow, slip flow, transitional flow, and
60 molecular-free flow regimes in micro and nano scales [5]. Later on, Civan (2010) developed

61 an equation to describe the behavior of permeability based on the Beskok and Karniadakis'
62 equation [7]. Javadpour et al. (2007 and 2009) introduced a model for the behavior of apparent
63 permeability of nanopores when Knudsen diffusion and slip flow are presented [8, 9].

64 The major portion of the gas in-place of shale gas reservoirs is found as free gas and or
65 adsorbed gas, while a small portion of it is dissolved in the formation water. Alongside with
66 different flow regimes, adsorption mechanism affects the performance of the shale gas
67 reservoirs by altering petrophysical properties of porous media, such as effective pore radius,
68 effective porosity, apparent permeability. Due to the importance of the adsorption, several
69 researchers have studied this phenomenon, where the surface diffusion is introduced as an
70 intrinsic property of gas by Fathi and Akkutlu (2009) and (2012) [10, 11]. The range of
71 1.55×10^{-7} - $8.80 \times 10^{-6} \text{ m}^2/\text{s}$ is reported by Kang et al. (2011) for the surface diffusivity of the
72 kerogen pore surface [12]. Furthermore, the Langmuir theorem is usually applied to capture
73 the effect of adsorption by using Langmuir-Freundlich (L-F) equation and modified Langmuir
74 (M-L) [13-15]. In another study, She et al. (2014) showed that apparent permeability would be
75 affected by adsorption significantly, especially in low-pressure conditions [16]. Wang et al.
76 (2016) showed that ignoring the changes in the petrophysical properties of porous media due
77 to adsorption/desorption cause an overestimation of original gas in-place and gas production
78 rate [17]. Zhang et al. (2018) proposed a permeability model with considering the Knudsen
79 diffusion, viscous flow, and the surface diffusion of adsorbed gas. They figured out that, at a
80 low-pressure condition, pore radius and temperature would affect the apparent permeability
81 [18]. Foroozesh et al. (2019) studied the effect of slip flow and gas adsorption on the apparent
82 permeability of a set of Eagle Ford and Barnett shale samples through a pore network model.
83 It was concluded that the effect of desorption is more significant compared to slip flow on
84 altering the permeability. Additionally, it was shown that the slip effect was able to reduce the

85 effect of desorption due to a reduction in pressure drop upon activating slip flow in their
86 model [19].

87 In addition to the mentioned active mechanisms during gas flow, the governed in-situ
88 stresses play an important role in the shale gas reservoirs. Porous material can be described as
89 a solid structure with an internal structure consisting of connected and isolated void spaces that
90 can hold fluids. Reservoirs consist of this solid matrix as well as the hydrocarbon held within
91 the network of pores, both existing at a state of pressure equilibrium before production. As
92 extraction of hydrocarbon commences, the change in pressure equilibrium induces a transfer
93 of stress to the solid network causing the deformation of the solid matrix that can lead to a
94 reduction in pore spaces. The reduction in pore volume alters the pressure equilibrium once
95 more pushing out the hydrocarbons within the pores creating an iterating loop of events. This
96 brings us to the law of effective stress which states that this coupled behavior is a function of
97 applied stress and pore pressure [20]. The deformations that occurred within the solid frame
98 are due to effective stress. The dynamic interaction between fluid flow and stress can be
99 described using a set of governing equations formulated based on Biot's poroelasticity theory
100 [21]. Induced deformations are described through a set of constitutive equations of linear
101 elasticity.

102 Many present studies have taken stress effects into consideration while assessing the
103 permeability of either conventional or unconventional reservoirs [22-25]. This is usually
104 through experimental tests conducted on shale core samples to correlate effective stress and
105 permeability. Various logarithmic and empirical equations [26, 27] as well as exponential
106 functions [28-31] have been used to describe the relationship. In a work done by Yang et al.
107 (2017), Klinkenberg's model was extended to a stress-dependent permeability model. The

108 effects of confining pressure on permeability reduction were noted by Soeder (1988) [32]. A
109 study by Gutierrez et al. (2015) also concluded that a large reduction in fracture permeability
110 occurs as effective normal stress is increased [33]. Bustin et al. (2008) estimated the
111 permeability of Barnett, Muskwa, and Ohio shale core samples under confining pressures
112 ranging up to 40 MPa. They found an exponential relationship between the shale permeability
113 and effective stress [34]. Similarly, Dong et al. (2010) carried out several experimental tests on
114 Chinsui shale core samples located in Taiwan under a confining pressure range up to 120 MPa.
115 Although an exponential relationship was shown at lower confining pressures (up to 40 MPa),
116 at higher confining pressures this behavior was no longer maintained as it was more adequately
117 fitted with the power law. This shift in behavior could be due to the mechanical properties of
118 shale rocks [23]. Wentong et al. (2018) conducted a study on the stress sensitivity of shale
119 reservoirs. Their findings showed that pore compressibility shares a negatively linear
120 relationship with the rocks Young's modulus. It was observed that, at higher Young's modulus
121 values, pore compressibility would decrease. [35].

122 In addition, a few studies have been done on the modeling aspects of the permeability stress-
123 dependency in shale gas reservoirs. Nazari Moghaddam and Jamiolahmady (2016) studied the
124 effect of stress and slippage on the permeability of shale gas reservoirs. At constant effective
125 stress, a permeability enhancement was reported when pressure decreased due to gas slippage
126 phenomena [36]. Sun et al. (2017) developed a model for apparent permeability of the
127 nanotubes based on the Beskok and Karniadakis' equation, which shows a significant effect of
128 stress on the transport capacity [37]. Feng et al. (2019) showed one order of magnitude
129 reduction of the apparent permeability due to the effect of stress when slip flow, Knudsen flow,
130 and surface diffusion are contributing to a nanoscale model [38].

131 When it comes to pore-scale simulation, pore network modeling provides an effective way
132 to simulate the different physical and chemical phenomena, such as deep bed filtration, non-
133 Darcy flow, multiphase flow, reactive flow, and non-Newtonian displacement [39, 40]. Fatt
134 (1956) used pore network modeling for simulation of the capillary pressure curve as a function
135 of saturation for the first time [41]. After that, pore network modeling gets high attention in
136 different fields by researchers. For example, many researchers have tried to study fluid
137 dynamics by considering the effect of slip flow and Knudsen diffusion for the shale gas
138 reservoirs [8, 42-44]. The Hagen-Poiseuille equation is usually used in these studies for
139 simplification and saving computational cost, where it can reduce the accuracy of results.
140 Nevertheless, the Navier-Stokes equation has also been used instead of the Hagen-Poiseuille
141 equation by scholars [19].

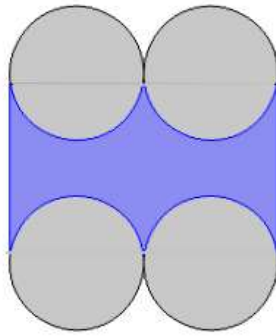
142 To the best of the authors' knowledge, a detailed pore network modeling study with considering
143 the slip flow regime, desorption of gas, and the effect of stress on the shale gas permeability,
144 has not been presented in the literature. Therefore, the objective of this study is to use pore-
145 scale models to investigate the contribution of slip flow and adsorption/desorption to shale gas
146 permeability while considering stress effects. In other words, this study aims to investigate the
147 apparent permeability of the shale rocks at different effective stress values when slip flow and
148 gas desorption are present. This study helps to have a better and realistic understanding of gas
149 flow behavior in shale rocks. Considering above-mentioned physics simultaneously applied on
150 the heterogenous network of pores and flow channels in the model, it provides a more accurate
151 tool for sound management of gas production from shale reservoirs.

152 To do so, a single pore and a pore network models were established using the pore structure
153 image of a Berea sandstone sample. The unsteady-state Stokes equation was used rather than

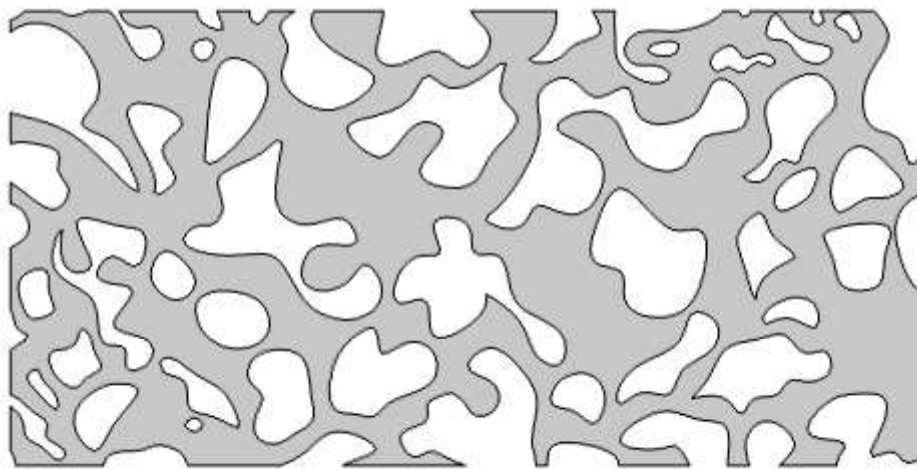
154 the Hagen-Poiseuille equation to capture the effect of fluid compressibility. Slip flow and gas
155 adsorption were included in the models by adapting appropriate boundary conditions. To
156 include the effect of stress, the solid deformation physics was included through the
157 poroelasticity term. Coupling fluid flow and stress will help to have a solid understanding of
158 apparent permeability alteration in shale gas reservoirs.

159 **2. Methodology**

160 Although the shale gas reservoirs have pores and throats size in the nanoscale, various ranges
161 are reported for the size of pores and throats by researchers. For example, Cipolla et al. (2009)
162 reported 1-200 nm as a typical range for pore diameters, while the reported size of the pore
163 diameters by Sakhaee-Pour and Bryant (2012) was 5 to 20 nm [45, 46]. Pore throats are
164 considered as connecting throats with the main resistance to flow. The average size of the pore
165 throats is in the range of 1-10 nm. Therefore, considering a valid range of size for the pores
166 and pore throats plays an important role in the simulation of the fluid flow in porous media.
167 This study is initially conducted on a single pore to quantify the effects of the physical
168 phenomena and flow mechanisms without the complexities of pore geometry. The geometry
169 of the single pore consists of four grains with a minimum pore throat diameter of 5 nm and a
170 length of 24 nm (Figure 1). Then, it is extended to a pore network by using an image of a Brea
171 sandstone pore structure when it is down scaled to the nanoscale (Figure 2). The dimension of
172 the pore network model in the x-direction and the y-direction is 705 nm and 352 nm,
173 respectively. In addition, the diameters of the pores are within the range of 5-20 nm.

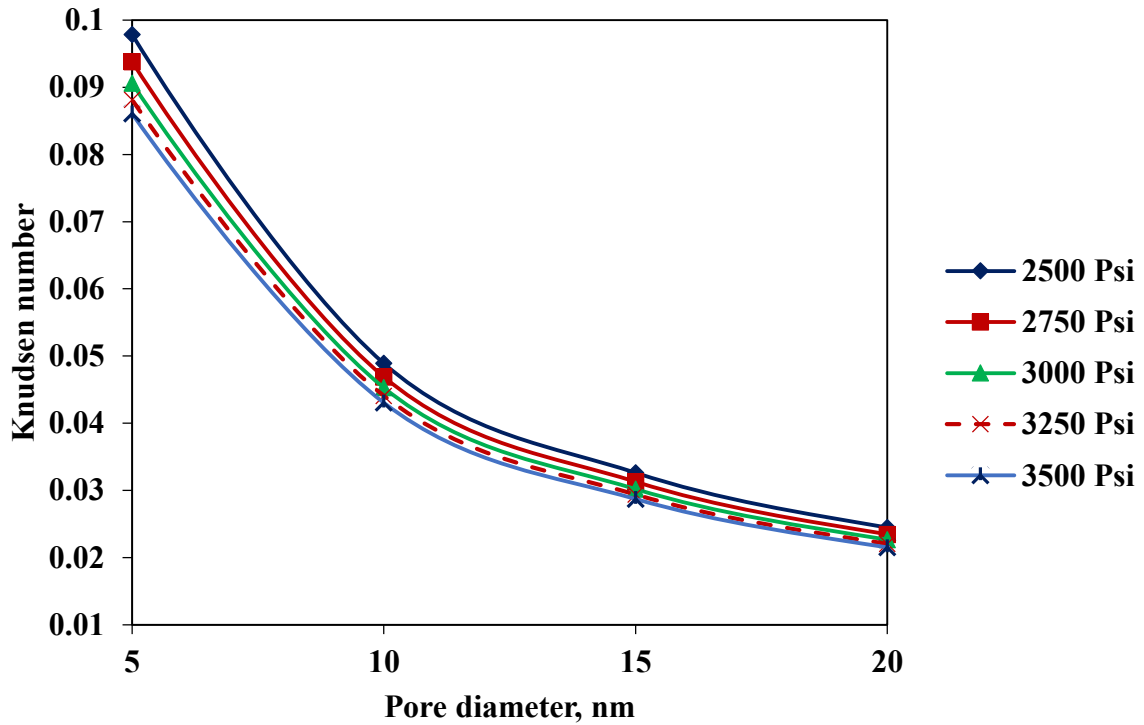


174
175 Figure 1. The geometry of the single pore model. The pathway of the fluid flow is highlighted
176 by the blue color.



177
178 Figure 2. The constructed pore network model. The white and the grey colors represent the
179 grains and the pores/throats, respectively [47]

180 When the size of the pores and pore throats of a porous media is comparable to or smaller
181 than the mean free path of the flowing fluid, the Knudsen diffusion is getting attention. To
182 ensure that the flow regime remains in the range of slip flow, the Knudsen number values are
183 calculated based on the applied pore pressure and size of the pores. There are two more benefits
184 to this issue. The first one is to exclude the Knudsen diffusion flow, and the second one is to
185 to maintain the validity of the Stokes flow. Figure 3 shows the Knudsen number with respect
186 to the pore diameter for different applied pore pressures. The results show the calculated
187 Knudsen numbers are in the range of continuum to slip flow.



188 Figure 3. Knudsen number with respect to pore diameters for different pore pressures
 189 embedded in the pore network studied for methane as a real gas.
 190

191 As mentioned, the flow regime is different at any given time of the life of a shale gas
 192 reservoir. It is governed by different parameters, such as reservoir pressure, geometrical aspects
 193 of the reservoir, and reservoir temperature. Based on the model setup applied in this work, the
 194 viscous flow with slip boundary conditions together with the desorption mechanism and stress
 195 effects are considered. In addition, the below assumptions are made in this study to establish a
 196 more reliable study and save the computational cost.

- 197 • The pores are fully saturated by methane (single phase).
- 198 • The methane is compressible, and its density and viscosity vary with pressure.
- 199 • The temperature is constant (40 °C).
- 200 • The gravity force is ignored (horizontal system).

201 As mentioned before, viscous flow plays a critical role to capture the fluid flow. To do so,
 202 the Navier-Stokes equation is used for the fluid flow of a viscous compressible Newtonian
 203 fluid, as presented in Eq. 1 [48].

$$\rho \left(\frac{d\vec{u}}{dt} + \vec{u} \cdot \nabla \vec{u} \right) = -\nabla p + \mu \nabla^2 \vec{u} + \frac{1}{3} \mu \nabla (\nabla \cdot u) \quad (1)$$

204 where ρ presents the density of the fluid in kg/m^3 , \vec{u} stands for the velocity field, p presents
 205 the pressure in Pa, and μ presents the gas viscosity in Pa.s.

206 The inertial force is negligible for the presented system of pores and throats because of its
 207 nanoscale. Therefore, the Navier-Stokes equation in the form of Eq.1 can be rewritten in the
 208 form of Eq.2 (Stokes equation) when the inertial force is negligible [49].

$$\rho \left(\frac{d\vec{u}}{dt} \right) = -\nabla p + \mu \nabla^2 \vec{u} + \frac{1}{3} \mu \nabla (\nabla \cdot u) \quad (2)$$

209 The terms density (ρ) and viscosity (μ) in Eq.2 are varied with pressure. Eq.3 and Eq.4
 210 present the variation of density and viscosity with pressure, which are obtained based on the
 211 non-linear regression of NIST chemistry webbook data [50]. It is worth mentioning that Eq.3
 212 and Eq.4 are valid within the range of 1 MPa to 35 MPa and a fixed temperature of 40 °C.

$$\rho = -3 \times 10^{-21} p^3 + 8 \times 10^{-14} p^2 + 7 \times 10^{-6} p - 1.8655 \quad (3)$$

$$\mu = -3 \times 10^{-28} p^3 + 2 \times 10^{-20} p^2 + 6 \times 10^{-14} p + 1 \times 10^{-5} \quad (4)$$

213 where p stands for pressure in Pa.

214 Slip flow is commonly presented in fluid flow in shale gas reservoirs. The slip velocity is
 215 defined as presented in Eq.5 [51, 52].

$$u_s = L_s \left(\frac{\partial u}{\partial n} \right)_s \quad (5)$$

216 where u_s is slip velocity in m/s , L_s presents the slip length in m , and n stands for the
 217 coordinate normal to the wall. In addition, the slip velocity is also presented in the form of
 218 Eq.6.

$$u_s - u_w = C_1 \lambda \left(\frac{\partial u}{\partial n} \right)_s \quad (6)$$

219 where u_s is slip velocity, u_w presents the surface/wall velocity. It is considered to be zero in
 220 this study because there is no wall movement. λ stands for the gas mean free path (methane
 221 molecules) which is a function of pressure, and C_1 stands for the slip coefficient, which is
 222 dependent on the tangential momentum accommodation coefficient (σ_v) and expressed as Eq.7.
 223 The value of 0.9 is used for σ_v in this study [53].

$$C_1 = \alpha_s \frac{2 - \sigma_v}{\sigma_v} \quad (7)$$

224 In Eq.7, α_s is a constant which is considered equal to one in the Maxwell model [54].

225 The desorption affects the properties of porous media, such as porosity, permeability, and
 226 pore radius. This effect can be described by the Langmuir theorem [55]. Eq.8 presents the
 227 detailed form of the Langmuir equation when it is converted into a volume unit.

$$vg = \frac{p \times vl}{p + pl} \times \rho_r \times h \times A \quad (8)$$

228 where vg presents gas content from the measured isotherm in m^3 , vl stands for the Langmuir
 229 adsorbed gas in m^3/kg , ρ_r presents the rock density in kg/m^3 , A presents the wall/rock/grain
 230 surface area in m^2 and h presents the wall/rock/grain thickness in m .

231 The adapted form of the Langmuir equation in cooperation with the mass balance equation
 232 is applied to the pores' wall as a boundary condition (Eq.9).

$$u_d = -\frac{vl \times \rho_r \times h}{\rho} \left[\rho \frac{d\left(\frac{p}{p+pl}\right)}{dp} + \frac{p}{p+pl} \frac{d\rho}{dp} \right] \frac{dp}{dt} \quad (9)$$

233 Needed for the above equation, Eq.3 is used to calculate the $\frac{d\rho}{dp}$ term. The rock density and
 234 rock thickness are assumed to be $2.4 g/cm$ and 4 microns, respectively [56]. Langmuir volume
 235 and pressure are extracted from the Heller and Zoback (2014) study for an Eagle Ford shale
 236 sample [57].

237 Governing equations are used to couple force equilibrium equations with the conservation
 238 of fluid mass. The equations include stress or strain factor that accounts for the poroelastic
 239 continuity equations, as well as the inclusion of pore pressure or fluid increment parameters
 240 that represents the fluid mass within the poroelastic system. To include the effect of stress in
 241 the presented model, the structural equations that are shown below are used to describe
 242 poroelasticity in the porous model based on the classical elasticity theory, which is derived in
 243 terms of shear moduli:

$$G\nabla^2 q + \frac{G}{1-2\nu} \left[\frac{\partial^2 q}{\partial x^2} + \frac{\partial^2 r}{\partial x \partial y} \right] + F_x - \alpha \frac{\partial p}{\partial x} = 0 \quad (10)$$

$$G\nabla^2 r + \frac{G}{1-2\nu} \left[\frac{\partial^2 q}{\partial y \partial x} + \frac{\partial^2 r}{\partial y^2} \right] + F_y - \alpha \frac{\partial p}{\partial y} = 0 \quad (11)$$

244 In Eqs 10-11, G is the shear modulus, q and r are displacements of solid (rock) in x and y
 245 directions, and ν is the Poisson ratio, α is Biot's constant, p is pore pressure.

246 The mechanical model solved in this study is then given through the following equations:

$$-\nabla \cdot \sigma = F_v, \sigma = S \quad (12)$$

$$S - s_0 = C: (\varepsilon - \varepsilon_0 - \varepsilon_{inel}) - \alpha p I \quad (13)$$

247 The strain is expressed in its index notation as shown in Eq.14.

$$\varepsilon_{ij} = \frac{1}{2} \left[\frac{\partial q_i}{\partial x_j} + \frac{\partial q_j}{\partial x_i} \right] \quad (14)$$

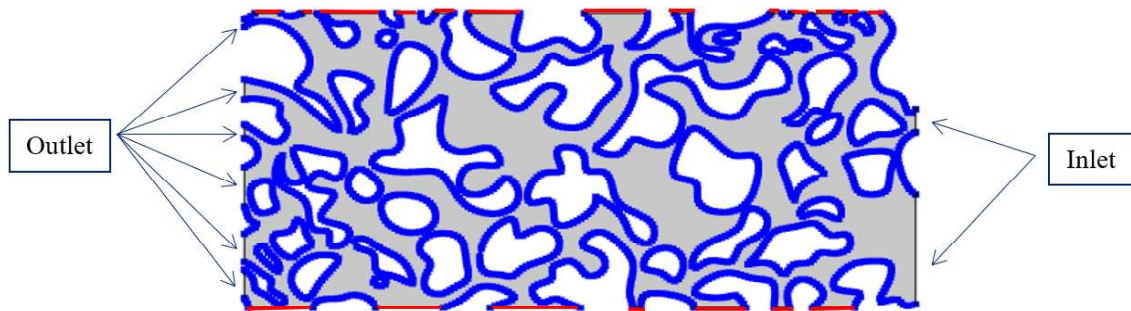
248 The stress tensors are denoted through σ and S , the volume force vector being expressed by
 249 F_v , with the strain tensor term being ε and deformation vector term being q . The pore pressure
 250 is denoted through the letter p , elastic tensor through the letter C and Biot's constant through
 251 the term α . Subscript '0' describes initial values, whereas 'inel' indicates the inelastic part of
 252 the strain.

253 In the last step, Darcy's equation is used to calculate the apparent permeability of the whole
 254 model, as presented in Eq.15.

$$K = \frac{V \times \Delta x \times \mu}{\Delta p} \quad (15)$$

255 where K presents apparent permeability of the porous media in D (Darcy), V presents the
 256 average velocity of gas in the pore network in cm/s , Δx stands for the length in the x-direction
 257 in cm, μ stands for viscosity in cp, and Δp is $(p_{inlet} - p_{outlet})$ in atm.

258 The above-mentioned equations were solved in an unsteady state condition using a finite
 259 element platform. To do so, the Stokes flow equation is solved by considering slip flow and
 260 gas adsorption as contributing terms in the boundary conditions. The inlet and outlet pressures
 261 of the pore network are assigned to be constant. Figure 4 shows the inlet (right side) and outlet
 262 (left side) boundaries as well as top and bottom boundaries (red lines). Pore walls are shown
 263 using blue lines where the slip flow and gas desorption are applied as boundary conditions
 264 using proper velocity term to the pore walls shown above. The pore pressure varies between
 265 2500-3500 psi based on reported values of the pressure for the shale reservoir reservoirs [56].
 266 The pressure gradient is set at 2 MPa/m (1.4 Pa/nm) [19]. The total simulation time is 2 seconds
 267 with a time step size of 0.1 seconds.



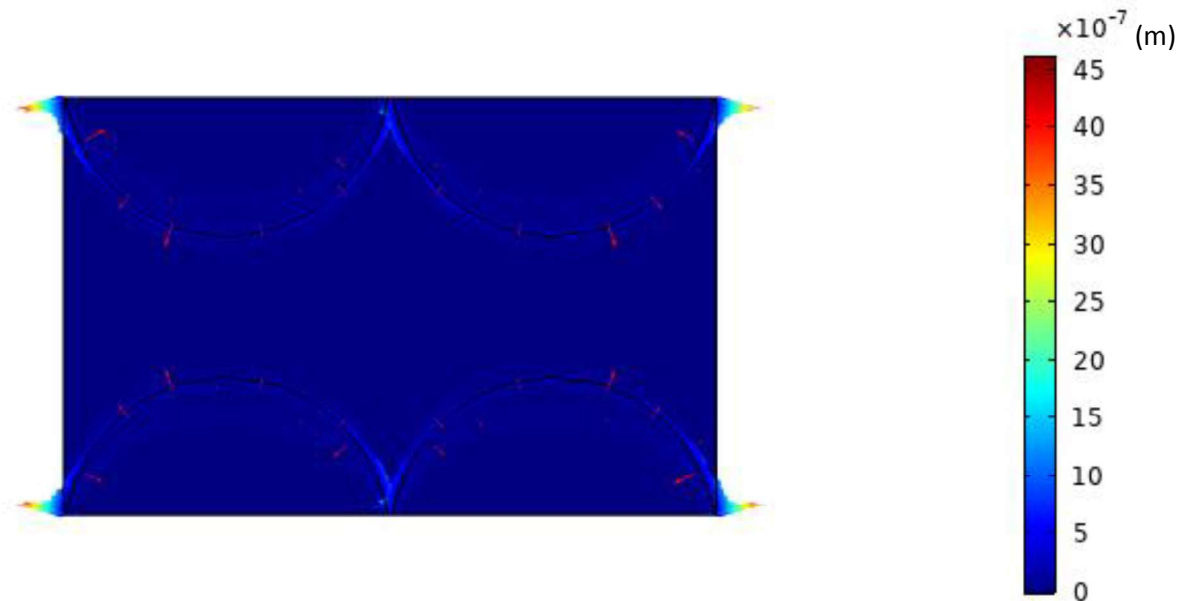
268
 269 Figure 4. Pore network model

270 3. Results and discussion

271 This section discusses the flow dynamics considering slip flow and gas desorption under
272 different effective stresses for a single pore model and a pore network model.

273 **3.1 Single pore model**

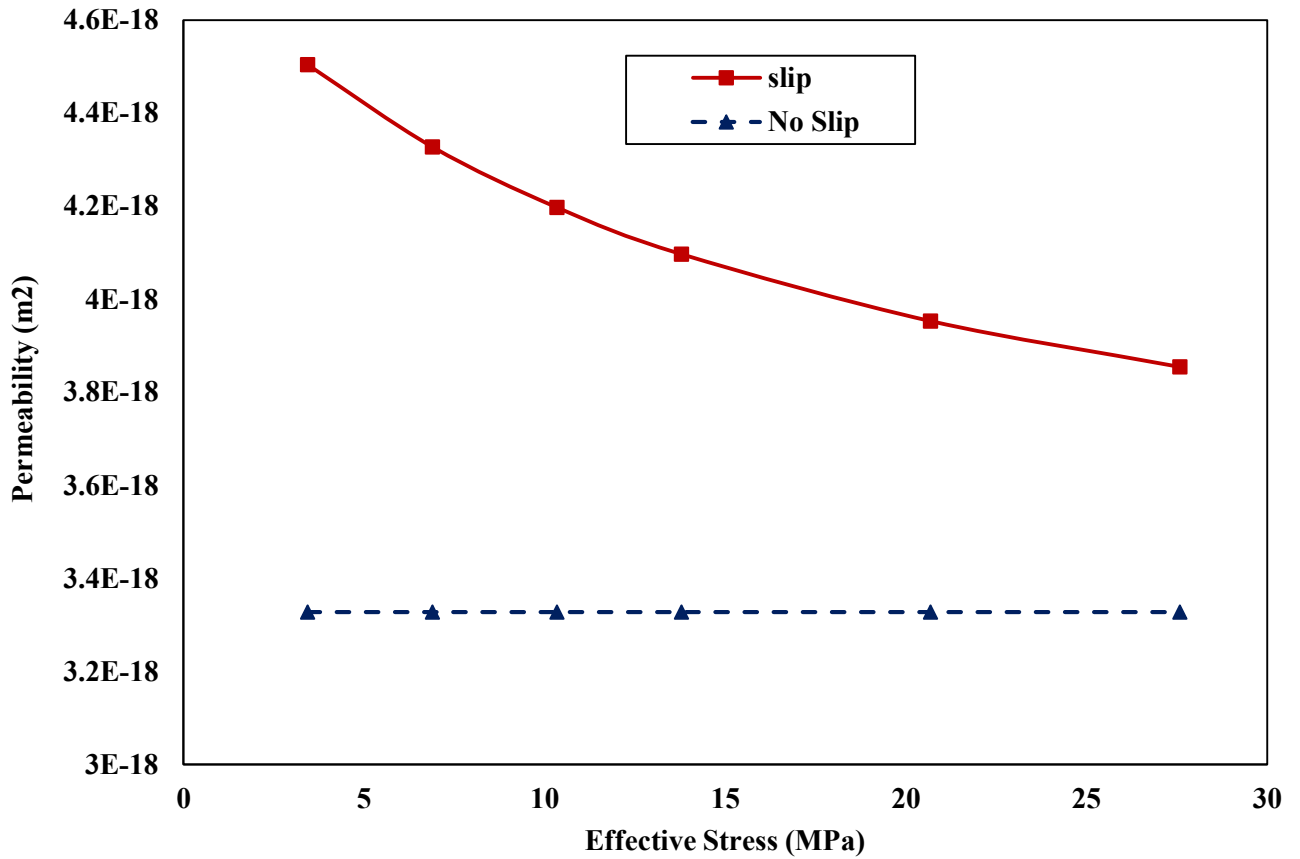
274 The single pore model was run under different effective stresses, where only the Stokes
275 equation was used to describe the fluid flow (no slip or desorption). This case is set as a
276 reference case to compare to cases that include the slip flow and gas desorption mechanisms
277 later. The model was run with constant pore pressure of 13.8 MPa (2000 psi), while the
278 effective stress was gradually increased from 3 MPa to 27 MPa (which is equivalent to 500
279 psi to 4000 psi) by increasing overburden pressure (loading stress). This range was chosen
280 based on effective stress data gathered from Barnett and Eagle Ford shale reservoirs, where
281 effective stress within Barnett shale reached levels up to 30 MPa [58]. Figure 5 shows the
282 deformation of the pores' wall for the single pore model at the end of the simulation time. A
283 deformation in pores' wall and contently in pore space could be noticed which was within the
284 range of 4-10 nm. It is worth mentioning that, in this single pore system, the wall deformation
285 has occurred almost in all directions, but more deformation can be observed horizontally (see
286 Figure 5) causing the pore space change to be minimal.



287 Figure 5. Deformation of the single pore model. The arrows point out the direction of the
 288 pore's wall deformation.
 289

290 3.1.1 Effect of slip flow

291 The slip flow boundary condition is applied to the single pore model. Figure 6 compares
 292 permeability profiles at different effective stresses for two boundary conditions: 1) without the
 293 slip flow , and 2) with the slip flow .



294 Figure 6. Single pore permeability profile for no-slip flow and slip flow conditions at
 295 different effective stresses
 296

297 As shown in Figure 6, when effective stress increases the permeability values show
 298 insignificant changes for the no-slip flow condition. A 0.3% difference was observed by
 299 increasing effective stress from 3 MPa to 27 MPa. It is believed that this observation stems
 300 from the fact that the pore space change is minimal in this case as shown in Figure 5.

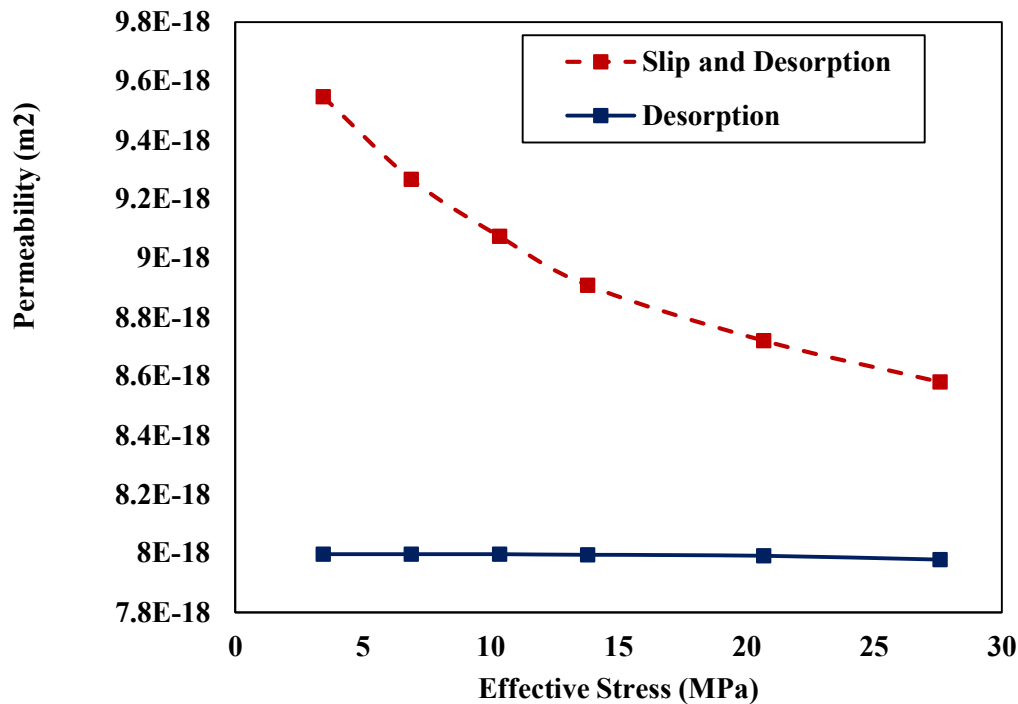
301 The permeability values are higher when the slip flow boundary is included. The slip flow
 302 leads to a 36% increase in permeability at the lowest-stress point compared to the no-slip case,
 303 while this value is 15% at the highest-stress condition. It should be noted that due to
 304 compression by increasing effective stress, the size of the pore reduces which is favorable for
 305 slip effect but at the same time, the pore pressure increases (due to pore space reduction) that

306 suppresses the slip effect. Considering the deformation is minimal here, pore pressure changes
307 would affect the slip flow physics more as compared to the pore size change. As a result, at
308 higher effective stress values, it is expected to have higher pore pressures (due to more
309 compression) and consequently have a lower slip flow contribution. This observation is in line
310 with the reported results in the literature, where it is mentioned that the effect of slip flow on
311 permeability is more dominated in lower pore pressure values [59, 60].

312 The permeability in the slip flow case is decreasing by 13% over increasing effective stress
313 to 27 MPa. The results show the slip flow mechanism is sensitive to stress because the slippage,
314 as a wall effect, can be influenced by wall deformation. It can also be noticed that with the
315 consideration of slip flow, the permeability trend shows an exponential behavior, where a
316 sharper trend of permeability reduction is observed at lower effective stresses due to higher
317 slip flow contribution.

318 **3.1.2 Coupled gas desorption and slip flow**

319 The gas desorption is captured by the model through a velocity boundary condition term as
320 mentioned before. Figure 7 shows the permeability profiles induced by the gas desorption and
321 coupled slip flow and gas desorption for the Eagle Ford shale sample.

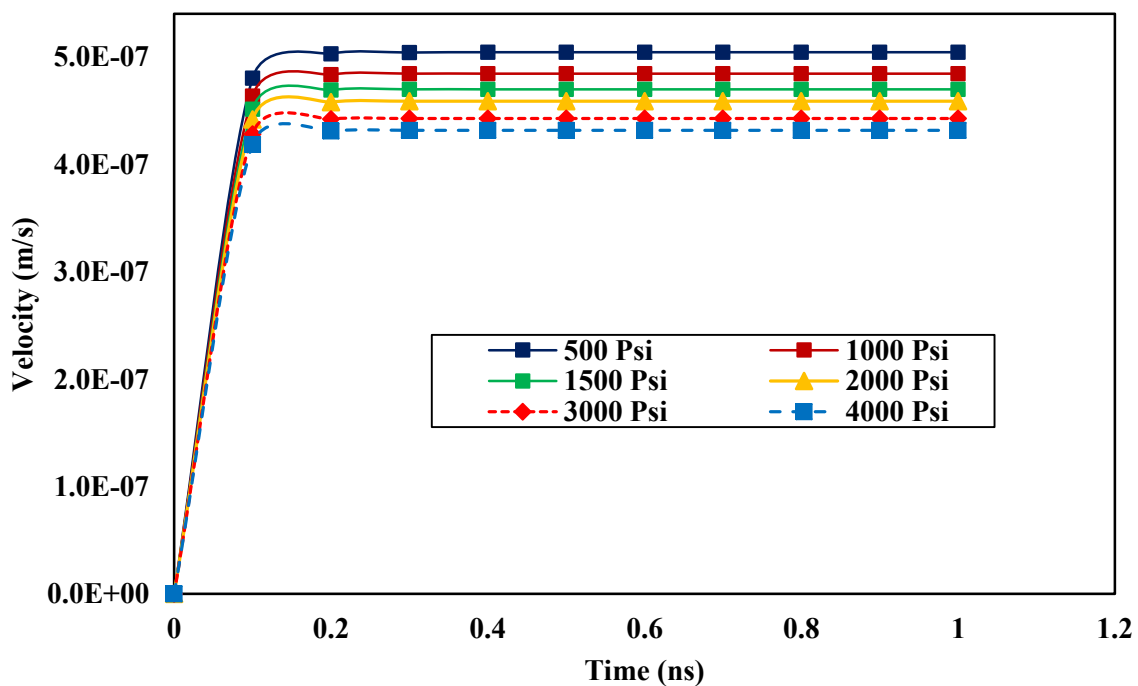


322
 323 Figure 7. Permeability profile developed by gas desorption and slip flow mechanisms for
 324 different effective stresses (single pore model).

325 Figure 7 shows the permeability profiles developed by the consideration of the gas
 326 desorption mechanism alone, as well as the coupled effect of gas desorption and slip flow
 327 mechanisms against effective stress. In the case of gas desorption only, a negligible
 328 permeability reduction is observed by increase of effective stress and the reductions has mainly
 329 happened at higher effective stress values. As mentioned earlier, the desorption mechanism is
 330 controlled by the pressure difference between Langmuir pressure and pore pressure.
 331 Considering the fact that an increase in the effective stress leads to an increase in pore pressure,
 332 hence it is limiting the desorption mechanism. Even more, the reduction in pore space can be
 333 attributed to two factors, 1) reduction as a result of rock matrix deformation and 2) a reversal
 334 of the desorption effect due to high pressure. Therefore, the increase in pore pressure could
 335 lead to the adsorption of methane which would act as a further reduction in pore diameter,
 336 resulting in a steeper decline in permeability at higher effective stresses.

337 In the scenario in which both slip flow and gas desorption mechanisms are activated, the
 338 permeability profile follows a similar trend to that developed by the slip flow mechanism in
 339 Figure 6. This is expected as when the effect of stress on individual desorption and also on
 340 individual slip flow contribution was discussed above, it was observed that that stress has no

341 significant effect on desorption mechanism individually while it has shown a considerable
 342 effect on individual slip flow contribution. Nevertheless, a higher permeability reduction is
 343 observed when the slip and desorption effects are coupled as compared to individual slip effect
 344 case (comparing Figure 6 and Figure 7). This is because at higher effective stress values,
 345 desorption contribution will be reduced resulting in a sharper decrease in the apparat
 346 permeability curve of slip-desorption (Figure 7) as compared to curve of slip alone (Figure
 347 6). Figures 8-10 show the velocity profile versus time for the three cases; slip flow acting alone,
 348 gas desorption acting alone and the coupled effect of slip flow and gas desorption to
 349 demonstrate the stress effect on these physics.

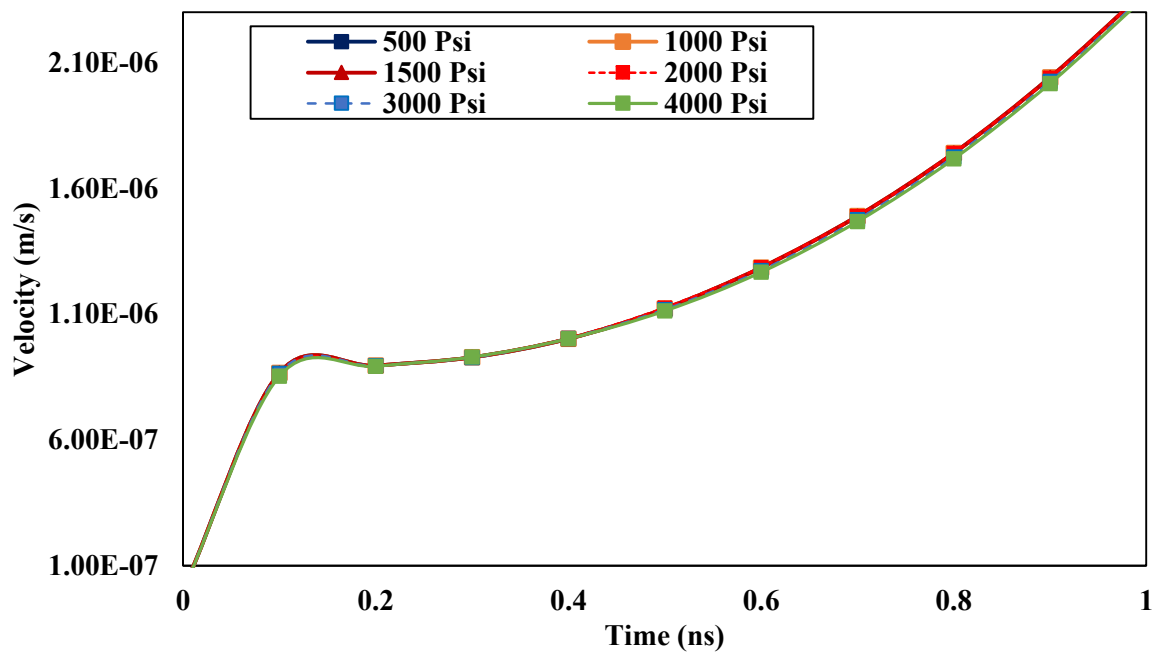


350 Figure 8. Velocity profile vs time under the effects of slip mechanism and stress effect
 351

352 When considering the stress effect on the slip mechanism, as shown in Figure 8, it can be
 353 noted that as effective stress increases, the flow velocity decreases due to deformation at higher
 354 effective stresses. Results imply that higher stresses cause more compression of the gas due to
 355 pore size reduction, resulting in pore pressure increment while the velocity decreases to satisfy
 356 conservation of energy law. It also can be observed that, at the early stages of flow, the velocity

357 is increasing before becoming steady for all the applied effective stresses that shows a transient
358 state due to gas compressibility at the beginning of the process.

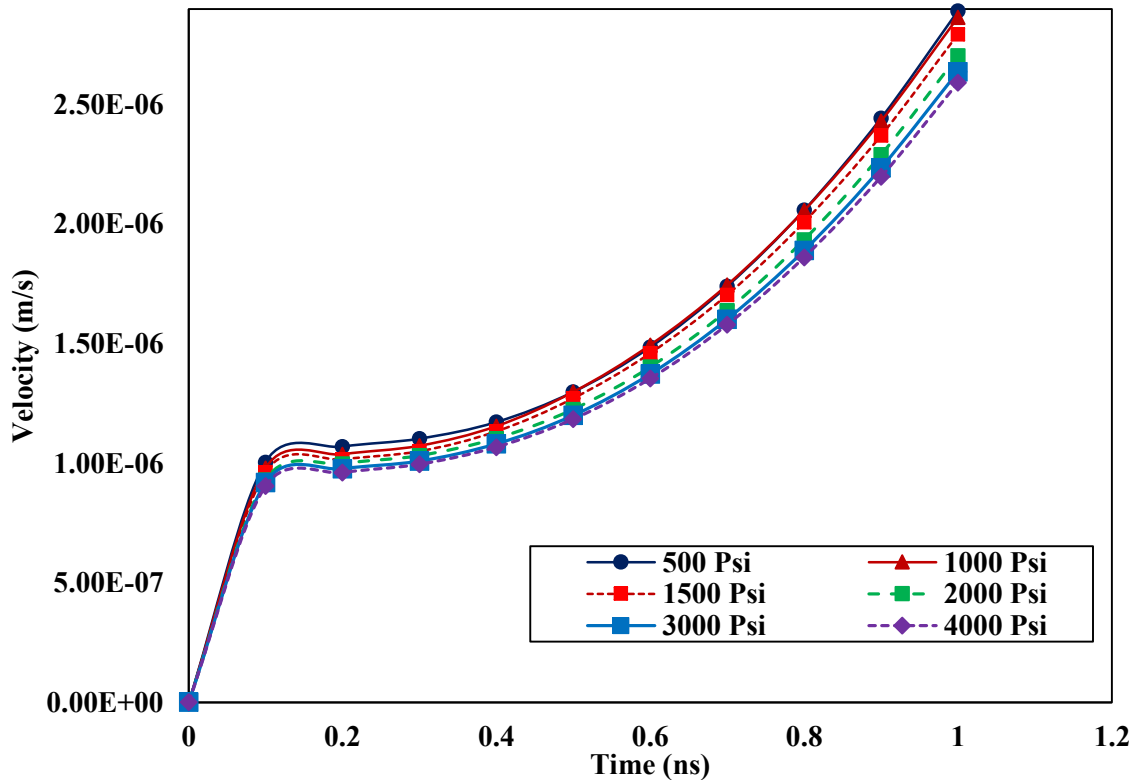
359 Figure 9 shows the velocity profile for the scenario in which gas desorption is evaluated at
360 different effective stresses showing insignificant effect of stress on desorption due to having
361 almost the same velocity profile for all the stress cases. The general trend of velocity is
362 increasing with time for all the effective stresses. A marginal reduction in velocity profiles can
363 be observed in the simulation time between 0.1 ns and 0.3 ns. This highlights the fact that
364 velocity has not reached a steady-state regime. Transient behavior of flow is due to the transient
365 nature of desorption mechanism.



366
367 Figure 9. Velocity profile vs time under the effects of gas desorption mechanism and stress
368 effect

369 Figure 10 shows the velocity profile when considering the combined effect of stress, slip
370 flow, and gas desorption. Comparing with the Figure 9, the effect of stress is more recognizable
371 in Figure 10 by looking at the curves that are not overlapping. As it can be observed, Figure 10
372 shows a higher velocity at lower stress values that is due to slip flow contribution (as also

373 observed in Figure 8). Additionally, Figure 10 illustrates that the velocity is increasing by time
 374 (a transient behavior) for all the applied stress values due to the contribution of the gas
 375 desorption (as also observed in Figure 9).

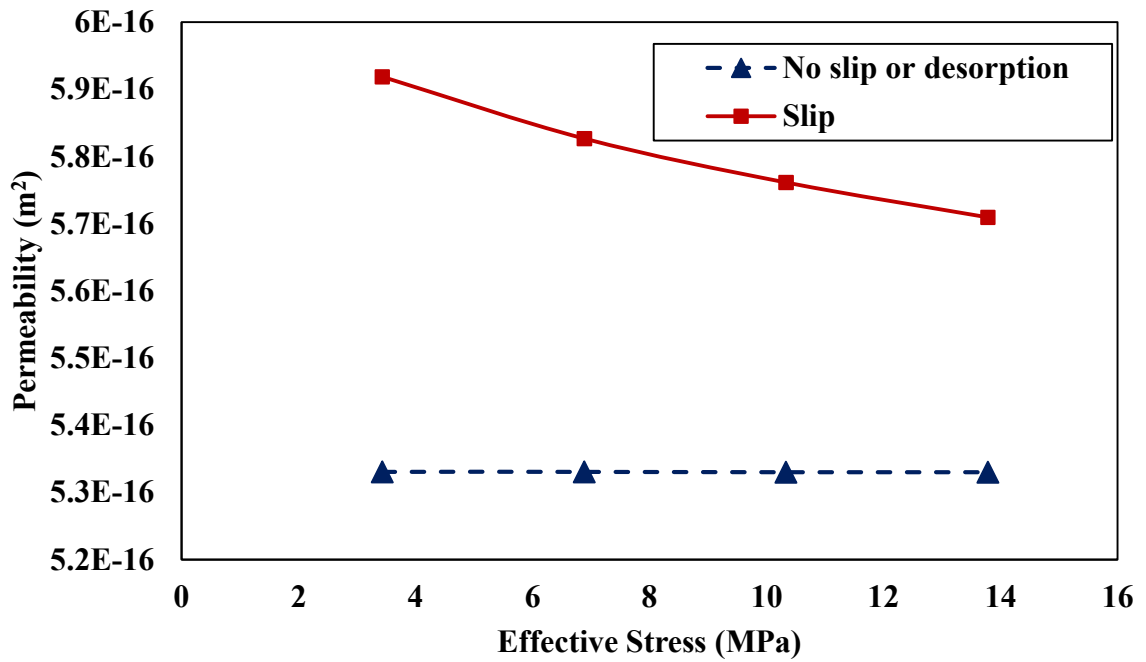


376 Figure 10. Velocity profile vs time under the coupled effects of slip flow, gas desorption, and
 377 stress effect
 378

379 3.2 Pore network model

380 3.2.1 Slip flow

381 Similarly, using a pore network, the stress effect was investigated. Figure 11 shows the
 382 results of the base case (no slip flow or gas desorption) against the case slip flow case. It is
 383 expected that at high effective stresses, the resulting deformation causes an intensive gas
 384 compression, making the slip flow contribution to become minimal. As a result, the two curves
 385 in Figure 11 are expected to cross at high-stress values.

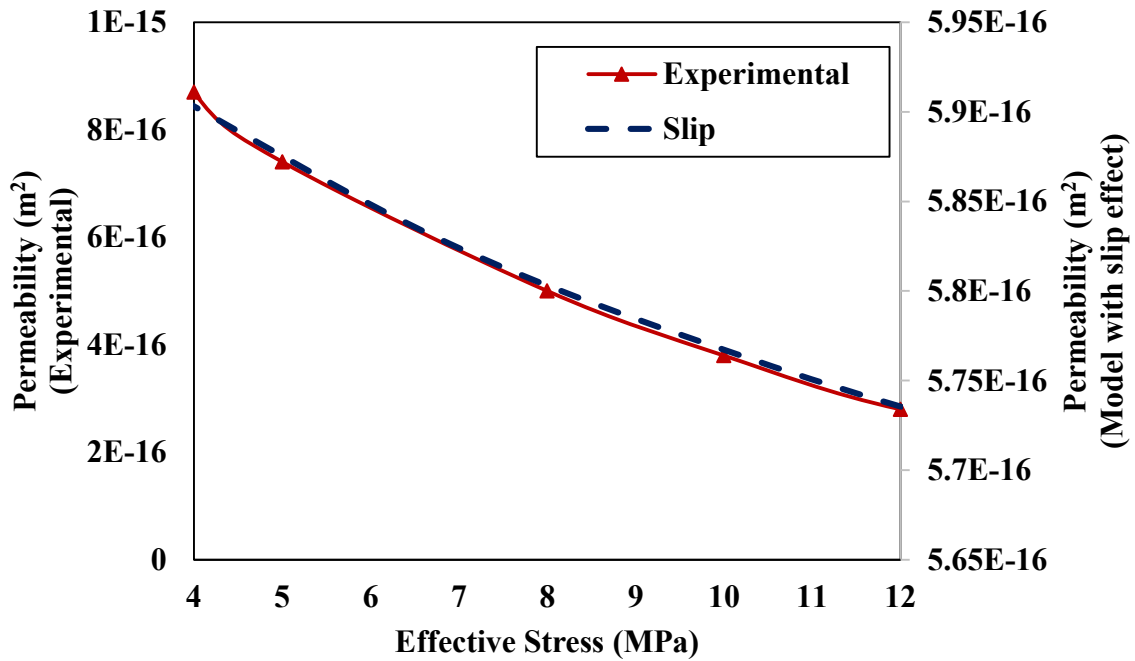


386
 387 Figure 11. Pore network permeability profile for no-slip flow condition and slip flow at
 388 different effective stresses

389 Figure 11 shows when fluid flow in the pore network is only affected by stress, the
 390 permeability remains constant. It also shows a reducing trend in the permeability of the pore
 391 network with respect to effective stress when the slip flow boundary is included. Increasing
 392 effective stress causes a reduction in pores and pore throats radius, results in more contribution
 393 of slip flow mechanism in fluid flow. A similar trend was observed when using the single pore
 394 model (see Figure 6)

395 To verify the model performance, the results produced by the model while considering the
 396 slip effect have been compared with experimental data gathered by Wei et al. (2018) on core
 397 samples received from Longmaxi shale formation [61]. Figure 12 shows that the trend of
 398 permeability profile predicted by the model is following the permeability trend obtained
 399 experimentally for a different type of shale sample. The experiment was conducted under
 400 constant injected gas pressure while loading pressure on the sample is gradually increased.

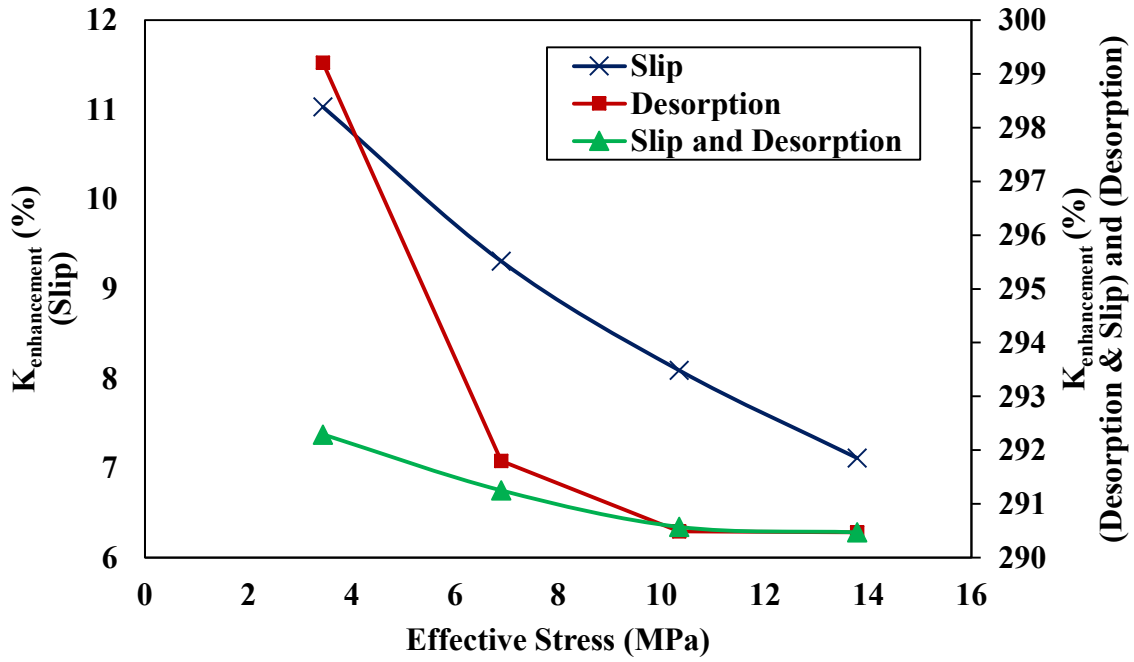
401 From the comparison, we can see that the expected behavioral trend obtained by our model
 402 follows the trend displayed by actual shale reservoirs. It should be noted that, as the shale
 403 sample used in the experiment is different from the sample checked by our model, it is not
 404 expected to get the same results. Nevertheless, predicting a similar trend with regard to stress
 405 effect is verifying that our model is set up properly.



406 Figure 12. Permeability profile prediction by the model with slip effect compared to
 407 experimental data
 408

409 3.2.2 Coupled gas desorption and slip flow effects

410 Figure 13 shows the permeability enhancement of a pore network due to slip flow, gas
 411 desorption, and coupled slip flow and gas desorption for an Eagle Ford shale sample.



412 Figure 13. Permeability enhancement versus effective stress
 413

414 Table 2 summarises the presented values of $K_{enhancement}$ (percent of permeability
 415 enhancement relative to the permeability of the base case) in Figure 13 for different
 416 mechanisms.

417 Table 2. Permeability alteration due to slip flow, gas desorption and coupled slip flow and
 418 gas desorption under various stress effects for the Eagle Ford sample

Effective stress (MPa)	$K_{enhancement}$ by slip flow (%)	$K_{enhancement}$ by gas desorption (%)	$K_{enhancement}$ by slip flow and gas desorption (%)
3.45	11.03	299.21	292.29
6.89	9.31	290.50	291.25
10.34	8.09	290.49	290.57
13.79	7.11	290.47	290.47

419 At the lowest effective stress applied, i.e., 3.45 MPa (equivalent to 500 psi), the
420 enhancement due to the slip flow mechanism is 11%, and as effective stress was gradually
421 increased the enhancement is reduced to 7% at 13.79 MPa (equivalent to 2000 psi). There are
422 two reasons for reducing trend of permeability enhancement. First, the diameter of flow
423 channels is reducing with increasing effective stress. The second, the rise of pore pressure acts
424 to limit the slip flow effect.

425 The scenario in which gas desorption is encountered showed that it is more sensitive to stress
426 within the pore network. This stress-sensitive behavior was not captured by single pore model
427 due to exclusion of pore size and structure heterogeneity in the single pore model. At 3.45
428 MPa (equivalent to 500 psi) effective stress, the enhancement is 299%. However, it has dropped
429 to 290% as the effective stress increases to 6.89 MPa (equivalent to 1000 psi), and as effective
430 stress keeps increasing from 6.89 MPa to 13.79 MPa, the reduction of permeability is
431 insignificant (only 0.03%). This could be due to the Langmuir pressure limit. As the pressure
432 in the system rises beyond the Langmuir pressure, the desorption mechanism is reversed, and
433 methane gas starts to adsorb to the surface. The Langmuir pressure for Eagle ford shale is 4.8
434 MPa (approximately 700 psi). Therefore, the increase in effective stress acts to limit the gas
435 desorption effect as 6.89 MPa (equivalent to 1000 psi) stress was applied. The negligible
436 permeability reduction observed in the range of 6.89 MPa to 13.79 MPa is due to the formation
437 of an adsorbed layer that acts to further reduce or restrict the flow diameter. This observation
438 is consistent with Wang and Marongiu-Porcu (2015) study where they reported that the effect
439 of gas adsorption/desorption is significant when effective stress is small such as near the
440 fracture area in a fractured formation as having smaller effective stress magnitude compared
441 to spots in the matrix zone [62]. In addition, it is reported in the study by Taghavinejad et al.

442 (2020) that permeability remains almost constant at higher pressures when the adsorption layer
443 is present [59].

444 The coupled effect of slip flow and gas desorption mechanisms show an enhancement of
445 292% at the applied 3.45 MPa (equivalent to 500 psi) stress and it falls to 290 % at 13.79 MPa
446 (equivalent to 2000 psi) stress. The slip flow and gas desorption act oppositely when both of
447 them are activated. It means that the magnitude of net permeability enhancement, which is due
448 to slip flow and gas desorption, is not equal to the summation of the magnitude of the
449 permeability enhancement due to individual slip flow and gas desorption. However, for higher
450 stress values, as the slip flow mechanism becomes more limited, we can notice that the curve
451 of coupled effect shows a closer matching to the curve of gas desorption alone. That is, slip
452 flow effect is no longer limiting the gas desorption, hence we can see a closer matching between
453 slip-desorption curve and desorption curve at higher stress values. In a study, Jia et al. (2020)
454 investigated the stress-dependent permeability values under slip flow and coupled slip flow
455 and gas desorption effects where a similar trend was observed verifying our results [60].

456 The inclusion of the stress term showed some notable results that help us understand the
457 nature of the stress effect on shale gas flow. Our model (both single and pore network model)
458 demonstrated that stress effect is limiting the gas flow in shale reservoirs by influencing both
459 slip flow and gas desorption contributions. It should be noted that the contributions of slip
460 flow and gas desorption are shown through an apparent permeability that represents the
461 enhancement to the permeability of the base case (the model without slip flow and desorption
462 terms). It was noted that the slip flow contribution would be limited by the stress effect. This
463 is because, although the reduction in pore volume due to deformation acts in favor of the slip
464 effect as gas-wall interaction increases, the resultant increase in pore pressure acts to limit the

465 slip effect as it restricts the mean free path of the flow. The sensitivity of the slip mechanism
466 to additional stresses was noted when a small incrementation in effective stress showed a
467 notable reduction in permeability enhancement. On the other hand, the effect of stress on the
468 desorption term was noted at higher effective stress values. The reduction in permeability
469 enhancement at low effective stresses was minor, this could be attributed to the gas adsorptive
470 nature of shale that is a function of the pore pressure. The increase in effective stresses leads
471 to an increase in pore pressure, which leads to the reversal of the desorption process into an
472 adsorption process. This in return would reduce the permeability enhancement due to reduction
473 of desorption contribution and furthermore would add to the reduction of pore space due to the
474 formation of an adsorbed layer on the wall surface.

475 **4. Conclusion**

476 In this study, the gas flow behavior of the shale gas reservoirs is investigated in pore-scale
477 in presence of slip flow, gas desorption, and the effect of stress. A typical single pore model
478 and a pore network models using a down-scaled Berea sandstone image were created. Although
479 the single pore model and pore network model produced almost similar trends of the mentioned
480 mechanisms, they were slightly different. The differences are attributed to the pore-scale size
481 and the additional heterogeneity found within the pore network model. Both models showed
482 negative effect of stress on gas flow when slip flow and desorption mechanisms were activated
483 in the models. In the single pore model, slip had a positive effect on desorption at a specific
484 stress value when their coupled contributions were summative. However, the results of pore
485 network model showed that slip had a destructive effect on desorption contribution specially
486 at low effective stress values. It should be noted that, the flow in the model is controlled by a
487 fixed differential pressure as constrained by the fixed inlet and outlet pressure boundary

488 conditions. As such, activating slip flow would result in having a higher pore pressure value
489 compared to the base case regardless of the effective stress value. Therefore, the changes in
490 pore pressure induced by the slip mechanism act to limit the desorption contribution. This was
491 clearly noted with the consideration of stress effects. At high effective stress values, the slip
492 effect is limited due to the increase in pore pressure by compression, hence we can notice a
493 closer matching between the permeability enhancement produced by the desorption mechanism
494 alone and the coupled effect of desorption and slip. Stress acts to limit the slip flow effect as
495 well as the desorption mechanism as they both are a function of pore pressure, however, as
496 stress increases the contribution of desorption is no longer reduced when both slip and
497 desorption are activated.

498

499

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